International Journal of Physical Research, 7 (1) (2019) 3-6



**International Journal of Physical Research** 

Website: www.sciencepubco.com/index.php/IJPR

Research paper



# Theoretical investigation of acoustic wave velocity of aluminum phosphide under pressure

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#### Abstract

The bulk and surface acoustic wave velocities of Aluminum phosphide (AIP) semiconducting material under pressure up to 9.5 GPa were studied. The structural parameters and the elastic constants used in this work are taken from our previous paper published in J. Optoelectron. Adv. M. 16, 207 (2014). The results obtained at zero-pressure are analyzed and compared with other data of the literature. In addition, the acoustic Grüneisen parameter and the Vickers and Knoop microhardness are predicted and analyzed in detail. Our calculated results are in good agreement with the experimental and other theoretical data of literature.

Keywords: Alp Semiconductor; Bulk and Surface Acoustic Wave Velocities; Newton-Raphson's Iterative Method; High Pressure.

# 1. Introduction

The binary compounds of III-V group's and their ternary alloys have been the subject of intense investigation in recent theoretical and experimental studies [1-7] due to their important physical properties such as: large bulk modulus, good hardness, and wide band-gap [1]. Among the binary compounds: AIP material, which crystallizes in cubic zincblende phase. It is an indirect wide band gap semiconductor (Eg = 2.48 eV), it has enormous technological importance in the industry of the infrared photo-detectors and light emitting diodes [1].

Recently, Aouadi et al. [2] studied some important physical properties of AlP material in both B3 and nickel-arsenide (B8) phases using PP-PW approach. Lakel et al. [3] preceded a systematic study of the effect of the pressure on the structural parameters, mechanical moduli and lattice dynamics properties of AlP material in both B3 and B8 phases.

Based on the idea, that the physical properties of elements and binary compounds can be improved artificially by adding some other constitutions to elaborate ternary alloys, Ameri et al. [4] have reported FP-LMTO calculations for equilibrium structure and the electronic properties of AlP and InP binary compounds and the effect of Al concentration on some physical properties of  $Al_xIn_{1-x}P$  ternary alloys combining AlP and InP binary compounds. The effect of the high pressure on the mechanical properties of BP, AlP, GaP and InP binary compounds were investigated by Bouhemadou et al. [5] using FP-APW + lo method.

Ehsanfar et al. [6] have performed a systematic investigation of the structural parameters, elastic constants, sound velocities, and some other physical properties of four III-phosphide binary compounds (BP, AIP, GaP and InP) in B3 phase using the DFT and the DFPT methods implemented in the Quantum ESPRESSO package. They also investigated the phonon spectrum of these compounds using the PHON program. They found that the frequency of the central optical phonon modes for AIP compound is around 427 cm<sup>-1</sup>, which is localized between 337 cm<sup>-1</sup> and 783 cm<sup>-1</sup> of GaP and BP, respectively [6].

In our previous work [7], the elastic and electronic properties of AlP material were investigated. In order to extend our research on this compound (AlP), we have calculated some other important parameters which are lacking in literature.

In the present work, as a part of this series of compounds, a bulk and surface-acoustic wave speeds up to 9.5 GPa are predicted. Our results obtained at equilibrium are analyzed and compared with the previous calculations and available experimental results of the literature.

# 2. Theory, results and discussion

## 2. 1. Elastic wave velocities for some specific crystallographic directions

Sound wave propagates through a solid by three acoustical modes, one of them is longitudinal and others two are transverse in nature. Hence, there are three types of velocities exits, one is the longitudinal wave velocity  $(V_p)$  and others two are the shear waves velocities  $(V_{sl}, V_{s2})$ . These three types of velocities depend on the direction of the propagation of the wave in the crystal [8]. The expressions for



direction dependent ultrasonic velocities in cubic crystals can be described as function of the elastic constants  $C_{ij}$  and the density  $\rho$  as follows [8]. The different relationships used here are detailed elsewhere [9]. Our obtained values of the longitudinal-wave mode speed  $(V_p)$  and transverse-wave mode speed  $V_{s1}$ ,  $V_{s2}$  in AIP material propagating in [100], [110] and [111] directions for different values of applied pressure up 9.5 GPa are given in Table 1, and compared with available data of the literature [6], [9], [10].

Cable 1: Sound Velocities in (m/s) for Major Directions in AIP Up to 9.5 GPa. <sup>a</sup> Longitudinal Waves, <sup>b</sup> Shea
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Propagation (Direction)	Plane of Polarization	At equilibrium ( $P = 0$ GPa)			
		This work	Ref. [6]	Ref. [9]	Ref. [10]
[100]	[100] <sup>a</sup>	7430	7710	7285	7970
	(100) <sup>b</sup>	5093	4550	5172	5090
[110]	[100] <sup>a</sup>	8235	7890	8138	8440
	(001) <sup>b</sup>	5093	4550	5172	5090
	(110) <sup>b</sup>	3649	4230	5213	4260
[111]	[111] <sup>a</sup>	8487	7950	8404	8590
	(111) <sup>b</sup>	4186	4340	4240	4560

Dran a setier (Direction)	Plane of Polarization 1.46		Pressure (GPa)			
Propagation (Direction)		1.46	3.17	4.87	6.50	9.50
[100]	[100] <sup>a</sup>	6730	6395	6187	6053	5892
	(100) <sup>b</sup>	4560	4274	4077	3933	3709
[110]	[100] <sup>a</sup>	7458	7083	6847	6691	6480
	(001) <sup>b</sup>	4560	4274	4077	3933	3709
	(110) <sup>b</sup>	3235	3000	2833	2833	2546
[111]	[111] <sup>a</sup>	7685	7298	7053	6891	6664
	(111) <sup>b</sup>	3730	3477	3300	3170	2985

As seen from this table,  $V_p$  are fastest along [111] direction and the shear waves are slowest along [100] direction. On comparing results of the sound velocities for major directions with those obtained by Ehsanfar et al. [6], Ma et al. [9] and Adachi [10], one can notice a good agreement, except for  $V_{s1}$  in the [110] direction, a slightly difference was observed. However, to the best of the authors' knowledge, no theoretical or experimental data of the sound velocities for major directions at high-pressure have appeared anywhere in literature to compare our results. Therefore, this work can be used as a reference for any further experimental or theoretical study of this compound.

## 2. 2. Elastic wave velocities of the aggregate material

For the aggregate materials (polycrystalline compounds), the Debye average velocity  $v_m$  is usually used to study the ultrasonic velocities wave propagation; this quantity is given by the following expression [11], [12]

$$v_{m} = \left[ \left( 1 / 3 \right) \left( \left( 2 / v_{1}^{3} \right) + \left( 1 / v_{1}^{3} \right) \right) \right]^{-1/3}$$
(1)

Where:  $v_1$  and  $v_t$  are the longitudinal and transverse elastic wave velocities respectively. They are determined by using the following expressions [11]:  $v_1 = ((3B + 4G)/3g)^{1/2}$ , and  $v_t = (G/g)^{1/2}$ , where: B is the bulk modulus, G is the shear modulus, and g is the density. The computed values of  $v_1$ ,  $v_t$  and  $v_m$  of AlP at zero-pressure are: 7994 m/s, 4455 m/s and 4960 m/s, respectively. They are specified in Table 2, and compared with available data of the literature [5], [9].

Our obtained results are in good agreement with the theoretical ones reported by Bouhemadou et al. [5] and Ma et al. [9] and a small deviation about 1.1 %. was found between our value (4960 m/s) of  $v_m$  and the value (5015 m/s) reported by Ma et al. [9].

Table 2: Sound Velocities (in m/s) of AIP Compound at Zero-Pressure, in Comparison with Other Data of the Literature [5], [9]

Parameter	Vl	VS	Vm
This work	7994	4455	4960
Ref. [5]	8426	4981	5518
Ref. [9]	7873	4514	5015

In Figure 1, the transverse, longitudinal, and average sound velocities are plotted against pressure. It is found that all the sound velocities decrease monotonically with increasing pressure.



Fig. 1: Longitudinal, Transverse and Average Sound Velocities of AlP Semiconductor versus Pressure.

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By best fitting the longitudinal, transverse, and average acoustic wave velocities, the following function relationships have been obtained  $(v_l, v_t \text{ and } v_m \text{ were expressed in km/s, and p in GPa})$ :

$$v_1 = 7.97 - 0.57 p + 8.02 \times 10^{-2} p^2 - 4.06 \times 10^{-3} p^3$$
(2)

There are various approximations to esteem the Grüneisen parameter  $\gamma$ , such as the thermodynamic  $\gamma$ , mode-specific  $\gamma$ , and acoustic  $\gamma$  [13]. This later quantity ( $\gamma$ ) depends on the longitudinal ( $v_1$ ) and transverse ( $v_t$ ) sound velocities, it can be estimated by the following expression [12], [13]

$$\gamma = (9/2) \left[ v_{i}^{2} - (4/3)v_{i}^{2} \right] / \left[ v_{i}^{2} + 2v_{i}^{2} \right]$$
(5)

At P = 0, the value of the Grüneisen parameter  $\gamma$  has been calculated and found to be 1.63; our obtained value of  $\gamma$  is slightly bigger than the value 1.49 reported by Adachi [10].

#### 2.3. Surface acoustic wave speeds

The speed of the sound plays a very important role in the surface acoustic wave (SAW) devices, filters, signal processing units, sensors and actuators [14], and widely used in different fields. Besides, the study of the elastic wave velocities for some specific crystallographic directions, the surface acoustic wave speeds of the AIP compound under high pressure have been calculated for the first time. The speed of the surface wave's Vsw in a B3 crystal propagating in the [100] and [110] directions are given as [15], [16]:

$$C_{11}\left(V_{w}^{2}-\frac{C_{4}}{g}\right)\left(V_{w}^{2}-\frac{C_{11}}{g}+\frac{C_{12}^{2}}{C_{11}g}\right)^{2}=C_{44}V_{w}^{4}\left(V_{w}^{2}-\frac{C_{11}}{g}\right)$$
(6)

$$C_{II}\left(v_{sw}^{2} - \frac{C_{44}}{g}\right)\left(v_{sw}^{2} - R + \frac{C_{I2}^{2}}{C_{IIg}}\right)^{2} = C_{44} v_{sw}^{4}\left(v_{sw}^{2} - R\right)$$
(7)

Where: R is expressed as follow:  $R = ((C_{11} + C_{12} + 2C_{44})/2g)$ . Equations (6) and (7) are of third degree in  $V_{sw}^2$  refer to the [100] and [110] directions, respectively. The pressure dependence of the surface acoustic wave speeds ( $V_{sw}$ ) propagating in the [100] and [110] directions is plotted in Figure 2.



It is observed that the surface acoustic wave speeds decreases with increasing pressure for both [100] and [110] directions. The solid lines in the curves represent the best fit of the calculated data and which can be approximated by the following analytical expressions respectively:

$$V_{sw} (kms^{-1}) = 4.088 - 0.329 \text{ p} + 4.17 \text{x} 10^{-2} \text{ p}^2 - 2.07 \text{x} 10^{-3} \text{ p}^3$$
(8)

$$V_{sw} (kms^{-1}) = 4.335 - 0.343 p + 4.45 x 10^{-2} p^2 - 2.25 x 10^{-3} p^3$$
(9)

It is important to note that the same behavior of the surface acoustic wave speeds under compression was also observed in cubic zincblende AlN compound [16] for both [100] and [110] directions. To the best of authors' knowledge, there are no data available in the literature on the surface acoustic wave speeds for AlP semiconducting material.

#### 2.4. Vickers and knoop microhardness

The mechanical behavior of solid can be specified by means of the microhardness, which can be calculated by various methods [15]. For the semiconductors, the Knoop and the Vickers indentations are usually used. The Vickers hardness  $H_V$  can be calculated by the following formula [1]:



$$H_{1} = 2 \left( G^{3} / B^{2} \right)^{0.585} - 3$$

Where G is the shear modulus, and B is the bulk modulus.

For some cubic semiconductors, the Knoop microhardness H and the Voigt shear modulus where Gv are related by [10]:

$$H = 0.139 \text{ Gv}$$

Our values of the Vickers hardness (Hv) and the Knoop microhardness H were found to be 6.11 and 6.96 GPa, respectively. The calculated values are respectively slightly smaller than the theoretical one (7.0 GPa) reported in Ref. [1], and higher than the experimental one (5.5 GPa) quoted by Adachi [10]. For some other cubic materials of group-IV, III-V and II-VI semiconductors, Adachi [10] proposed a relationship between the Knoop microhardness H and the inverse unit-cell volume ( $V^{-1} = a^{-3}$ ), and expressed as follow: H = (3220/V) - 14.6, where: V is in Å<sup>3</sup> and H is in GPa. Using the value of the lattice constant (a = 5.436 Å) [7], the Knoop microhardness H of AlP was found to be 5.45 GPa; which is in excellent agreement with the experimental one (5.5 GPa) quoted by Adachi [10], the deviation between the two values is only about 0.91%.

### **3.** Conclusion

Using the elastic constants and structural parameters obtained in our previous work published in J. Optoelectron. Adv. M. 16, 207 (2014), the bulk and surface acoustic wave velocities of Aluminum phosphide (AIP) under pressure up to 9.5 GPa were studied. The obtained results at zero-pressure are analyzed and compared with previous data of the literature. In general our findings are in good agreement with other data of the literature.

In addition, the Grüneisen parameter and the Vickers and Knoop microhardness are predicted and analyzed in details. Our obtained results of the microhardness are also in agreement with other data of the literature.

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