# He's semi-inverse method for Camassa-Holm equation and simplified modified Camassa-Holm equation 

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#### Abstract

In this Letter, we study Camassa-Holm equation and Simplified Modified Camassa-Holm equation by using the well-known He's Semi-inverse Method. The solitary solutions are obtained using the Ritz method. In fact, the He's Semi-inverse Method is a promising method to various systems of linear and nonlinear equations.


Keywords: Hes Semi-inverse Method, Solitary solution, Soliton equation.

## 1 Introduction

Many important phenomena and dynamic processes in physics, mechanics, chemistry and biology can be represented by nonlinear partial differential equations. The study of exact solutions of nonlinear evolution equations plays an important role in soliton theory and explicit formulas of nonlinear partial differential equations play an essential role in the nonlinear science. Also, the explicit formulas may provide physical information and help us to understand the mechanism of related physical models.
The study of exact solutions of nonlinear partial differential equations plays an important role in soliton theory and explicit formulas of nonlinear partial differential equations play an essential role in the nonlinear science. Also, the explicit formulas may provide physical information and help us to understand the mechanism of related physical models. Recently, many kinds of powerful methods have been proposed to find exact solutions of nonlinear partial differential equations, e.g., the tanh-method [1], the homogeneous balance method [2], homotopy analysis method $[3,4,5,6,7,8,9]$, the $F$-expansion method [10], three-wave method [11, 12, 13], extended homoclinic test approach $[14,15,16]$, the $\left(\frac{G^{\prime}}{G}\right)$-expansion method [17] and the exp-function method [18, 19, 20, 21, 22].
The outline of the present paper is as follows. In Sec. 2, we have a brief review on Camassa and Holm Equations.we introduce the He's Semi-inverse Method in Sec. 3. In Secs. 4 and 5 we apply the He's Semi-inverse Method on the Camassa-Holm equation and Simplified Modified Camassa-Holm equation, respectively. In those sections we obtain new and exact solutions for our equations. The paper is concluded in Sec. 6.

## 2 Camassa and Holm equations

Camassa and Holm [23] derived a completely integrable wave equation ( CH equation) for water waves

$$
\begin{equation*}
u_{t}+2 \alpha u_{x}-u_{x x t}+b u u_{x}=2 u_{x} u_{x x}+u u_{x x x}, \tag{1}
\end{equation*}
$$

by retaining two terms that are usually neglected in the small amplitude, shallow water limit. Tian and Song [24] investigated a modified CamassaHolm equation (MCH equation)

$$
\begin{equation*}
u_{t}+2 \alpha u_{x}-u_{x x t}+b u^{n} u_{x}=2 u_{x} u_{x x}+u u_{x x x}, \tag{2}
\end{equation*}
$$

and obtained new peaked solitary wave solutions. In addition, Boyd [25] investigated that if the solitary wave varies slowly with $\xi=x-c t$, then the two extra terms on the right-hand side of (1) will be small and the soliton is given
to lowest order by the solutions of

$$
\begin{equation*}
u_{t}+2 \alpha u_{x}-u_{x x t}+b u u_{x}=0 . \tag{3}
\end{equation*}
$$

In view of (3), Wazwaz [26] investigated a modified form of Camassa-Holm equation, which is simplified from MCH equation and given by

$$
\begin{equation*}
u_{t}+2 \alpha u_{x}-u_{x x t}+b u^{n} u_{x}=0 \tag{4}
\end{equation*}
$$

In this paper, we only consider $n=2$,

$$
\begin{equation*}
u_{t}+2 \alpha u_{x}-u_{x x t}+b u^{2} u_{x}=0 \tag{5}
\end{equation*}
$$

and for simplicity we call (5) simplified MCH equation. (For more details see [27])

## 3 Description of He's semi-inverse method

We suppose that the given nonlinear partial differential equation for to be in the form

$$
\begin{equation*}
P\left(u, u_{x}, u_{t}, u_{x x}, u_{t t}, u_{x t}, \ldots\right)=0 \tag{6}
\end{equation*}
$$

where $P$ is a polynomial in its arguments. The essence of He's semi-inverse method can be presented in the following steps:

Step 1. Seek solitary wave solutions of (6) by taking $u(x, t)=U(\xi), \xi=x-c t$ and transform (6) to the ordinary differential equation

$$
\begin{equation*}
U\left(u, u^{\prime}, u^{\prime \prime}, \ldots\right)=0 \tag{7}
\end{equation*}
$$

where prime denotes the derivative with respect to $\xi$.
Step 2. If possible, integrate (7) term by term one or more times. This yields constant(s) of integration. For simplicity, the integration constant(s) can be set to zero.

Step 3. According to Hes semi-inverse method, we construct the following trial-functional

$$
\begin{equation*}
J(U)=\int L d \xi \tag{8}
\end{equation*}
$$

where $L$ is an unknown function of $U$ and its derivatives.
There exist alternative approaches to the construction of the trial-functionals, see Refs. [28, 29].
Step 4. By the Ritz method, we can obtain different forms of solitary wave solutions, in the form

$$
\begin{equation*}
U(\xi)=p \operatorname{sech}^{n}(q \xi) \tag{9}
\end{equation*}
$$

where $P$ and $q$ are constants to be further determined.
Substituting (9) into (8) and making $J$ stationary with respect to $P$ and $q$ results in

$$
\begin{align*}
& \frac{\partial J}{\partial p}=0  \tag{10}\\
& \frac{\partial J}{\partial q}=0 . \tag{11}
\end{align*}
$$

Solving simultaneously (10) and (11) we obtain and. Hence, the solitary wave solution (9) is well determined.

## 4 New and exact solutions to the Camassa-Holm equation

In order to seek exact solutions of Camassa-Holm equation,

$$
\begin{equation*}
u_{t}+2 \alpha u_{x}-u_{x x t}+b u u_{x}=0 \tag{12}
\end{equation*}
$$

We suppose

$$
\begin{equation*}
u(x, t)=u(\xi) \quad, \quad \xi=x-w t \tag{13}
\end{equation*}
$$

where $w$ is complex constants to be determined later. Substituting (13) into (12), we have

$$
\begin{equation*}
-w u^{\prime}+2 \alpha u^{\prime}+w u^{\prime \prime \prime}+b u u^{\prime}=0 \tag{14}
\end{equation*}
$$

or equivalently

$$
\begin{equation*}
(2 \alpha-w) u^{\prime}+w u^{\prime \prime \prime}+b u u^{\prime}=0 \tag{15}
\end{equation*}
$$

where prime denotes the differential with respect to $\xi$. Integrating (15) with respect to $\xi$ and taking the integration constant as zero yields

$$
\begin{equation*}
(2 \alpha-w) u+w u^{\prime \prime}+\frac{b}{2} u^{2}=0 \tag{16}
\end{equation*}
$$

According to Ref. [28], By He's semi-inverse method [29], we can arrive at the following variational formulation:

$$
\begin{equation*}
J(\phi)=\int_{0}^{\infty}\left[-\frac{w}{2}\left(u^{\prime}\right)^{2}+\frac{(2 \alpha-w)}{2} u^{2}+\frac{b}{6} u^{3}\right] d \xi \tag{17}
\end{equation*}
$$

We assume the soliton solution in the following form

$$
\begin{equation*}
\phi(\xi)=p \operatorname{sech}^{2}(q \xi) \tag{18}
\end{equation*}
$$

where $p, q$ are an unknown constant to be further determined.
By Substituting (18) into (17) we obtain

$$
\begin{align*}
J & =\int_{0}^{\infty}\left(2 p^{2} w q^{2}+\frac{1}{6} b p^{3}\right) \operatorname{sech}^{6}(q \xi) d \xi+\int_{0}^{\infty}\left(-2 p^{2} w q^{2}+p^{2} \alpha-\frac{1}{2} w p^{2}\right) \operatorname{sech}^{4}(q \xi) d \xi \\
& =\frac{1}{q}\left(2 p^{2} w q^{2}+\frac{1}{6} b p^{3}\right) \int_{0}^{\infty} \operatorname{sech}^{6}(\theta) d \theta+\frac{1}{q}\left(-2 p^{2} w q^{2}+p^{2} \alpha-\frac{1}{2} w p^{2}\right) \int_{0}^{\infty} \operatorname{sech}^{4}(\theta) d \theta  \tag{19}\\
& =\frac{8}{15 q}\left(2 p^{2} w q^{2}+\frac{1}{6} b p^{3}\right)+\frac{2}{3 q}\left(-2 p^{2} w q^{2}+p^{2} \alpha-\frac{1}{2} w p^{2}\right),
\end{align*}
$$

For making $J$ stationary with respect to $p$ and $q$ results in

$$
\begin{align*}
& \frac{\partial J}{\partial p}=\frac{2 p}{15 q}\left(-4 w q^{2}+10 \alpha-5 w+2 b p\right)  \tag{20}\\
& \frac{\partial J}{\partial q}=-\frac{p^{2}}{45 q^{2}}\left(12 w q^{2}+30 \alpha-15 w+4 b p\right) \tag{21}
\end{align*}
$$

or simplifying

$$
\begin{align*}
& -4 w q^{2}+10 \alpha-5 w+2 b p=0  \tag{22}\\
& 12 w q^{2}+30 \alpha-15 w+4 b p=0 \tag{23}
\end{align*}
$$

From (22) and (23), we can easily obtain the following relations:

$$
\begin{equation*}
p=-\frac{3(w-2 \alpha)}{b} \quad, \quad q= \pm \frac{1}{2} \frac{\sqrt{w(w-2 \alpha)}}{w}, \tag{24}
\end{equation*}
$$

Therefore by (18), the solitary wave solutions can be approximated as

$$
\begin{equation*}
u(\xi)=-\frac{3(w-2 \alpha)}{b} \operatorname{sech}^{2}\left[ \pm \frac{1}{2} \frac{\sqrt{w(w-2 \alpha)}}{w}(x-w t)\right] \tag{25}
\end{equation*}
$$

for some arbitrary constants $w$.

## 5 New and exact solutions to the simplified MCH equation

In order to seek exact solutions of simplified MCH equation,

$$
\begin{equation*}
u_{t}+2 \alpha u_{x}-u_{x x t}+b u^{2} u_{x}=0 . \tag{26}
\end{equation*}
$$

We suppose

$$
\begin{equation*}
u(x, t)=u(\xi) \quad, \quad \xi=k x-w t \tag{27}
\end{equation*}
$$

where $w$ is complex constants to be determined later. Substituting (27) into (26), we have

$$
\begin{equation*}
-w u^{\prime}+2 \alpha u^{\prime}+w u^{\prime \prime \prime}+b u^{2} u^{\prime}=0, \tag{28}
\end{equation*}
$$

or equivalently

$$
\begin{equation*}
(2 \alpha-w) u^{\prime}+w u^{\prime \prime \prime}+b u^{2} u^{\prime}=0 \tag{29}
\end{equation*}
$$

where prime denotes the differential with respect to $\xi$. Integrating (29) with respect to $\xi$ and taking the integration constant as zero yields

$$
\begin{equation*}
(2 \alpha-w) u+w u^{\prime \prime}+\frac{b}{3} u^{3}=0 \tag{30}
\end{equation*}
$$

According to Ref. [28], By He's semi-inverse method [29], we can arrive at the following variational formulation:

$$
\begin{equation*}
J(\phi)=\int_{0}^{\infty}\left[-\frac{w}{2}\left(u^{\prime}\right)^{2}+\frac{(2 \alpha-w)}{2} u^{2}+\frac{b}{12} u^{4}\right] d \xi \tag{31}
\end{equation*}
$$

We assume the soliton solution in the following form

$$
\begin{equation*}
\phi(\xi)=p \operatorname{sech}(q \xi) \tag{32}
\end{equation*}
$$

where $p, q$ are an unknown constant to be further determined.
By Substituting (32) into (31) we obtain

$$
\begin{align*}
J & =\int_{0}^{\infty}\left(\frac{1}{2} p^{2} w q^{2}+\frac{1}{12} b p^{4}\right) \operatorname{sech}^{4}(q \xi) d \xi+\int_{0}^{\infty}\left(-\frac{1}{2} p^{2} w q^{2}+p^{2} \alpha-\frac{1}{2} w p^{2}\right) \operatorname{sech}^{2}(q \xi) d \xi \\
& \left.=\frac{1}{q}\left(\frac{1}{2} p^{2} w q^{2}+\frac{1}{12} b p^{4}\right) \int_{0}^{\infty} \operatorname{sech}^{4}(\theta) d \theta+\frac{1}{q}\left(-\frac{1}{2} p^{2} w q^{2}+p^{2} \alpha-\frac{1}{2} w p^{2}\right)\right) \int_{0}^{\infty} \operatorname{sech}^{2}(\theta) d \theta  \tag{33}\\
& =\frac{2}{3 q}\left(\frac{1}{2} p^{2} w q^{2}+\frac{1}{12} b p^{4}\right)+\frac{1}{q}\left(-\frac{1}{2} p^{2} w q^{2}+p^{2} \alpha-\frac{1}{2} w p^{2}\right),
\end{align*}
$$

For making $J$ stationary with respect to $p$ and $q$ results in

$$
\begin{align*}
& \frac{\partial J}{\partial p}=\frac{p}{9 q}\left(-3 w q^{2}+18 \alpha-9 w+2 b p^{2}\right)  \tag{34}\\
& \frac{\partial J}{\partial q}=-\frac{p^{2}}{18 q^{2}}\left(3 w q^{2}+18 \alpha-9 w+b p^{2}\right) \tag{35}
\end{align*}
$$

or simplifying

$$
\begin{align*}
& -3 w q^{2}+18 \alpha-9 w+2 b p^{2}=0,  \tag{36}\\
& 3 w q^{2}+18 \alpha-9 w+b p^{2}=0, \tag{37}
\end{align*}
$$

From (36) and (37), we can easily obtain the following relations:

$$
\begin{equation*}
p= \pm \frac{\sqrt{6 b(w-2 \alpha)}}{b} \quad, \quad q= \pm \frac{\sqrt{w(w-2 \alpha)}}{w} \tag{38}
\end{equation*}
$$

Therefore by (32), the solitary wave solutions can be approximated as

$$
\begin{equation*}
u(\xi)= \pm \frac{\sqrt{6 b(w-2 \alpha)}}{b} \operatorname{sech}\left[ \pm \frac{\sqrt{w(w-2 \alpha)}}{w}(x-w t)\right] \tag{39}
\end{equation*}
$$

for some arbitrary constants $w$.

## 6 Conclusion

In summary, by means of the modification of He's semi-inverse method, we obtained new and exact solutions for the Camassa Holm equation and Simplified Modified Camassa-Holm equation. He's semi-inverse method is a very dominant instrument to find the solitary solutions for various nonlinear equations.

## References

[1] A.M. Wazwaz, The tanh method: solitons and periodic solutions for the Dodd-Bullough-Tzikhailov and the Tzitzeica-Dodd-Bullough equations, Chaos, Solitons and Fractals, 25 (2005) 55-63.
[2] Z. Xiqiang, W. Limin, S. Weijun, The repeated homogeneous balance method and its applications to nonlinear partial differential equations, Chaos, Solitons and Fractals, 28(2) (2006) 448-453.
[3] S.J. Liao, Beyond Perturbation: Introduction to the homotopy analysis method, Chapman \& Hall/CRC Press, Boca Raton, 2003.
[4] S.J. Liao, On the homotopy analysis method for nonlinear problems, Appl. Math. Comput., 147 (2004) 499-513.
[5] S.J. Liao, A new branch of solutions of boundary-layer flows over an impermeable stretched plate, Int. J. Heat Mass Transfer, 48 (2005) 2529-2539.
[6] S.J. Liao, A general approach to get series solution of non-similarity boundary-layer flows, Commun. Nonlinear Sci. Numer. Simul., 14(5) (2009) 2144-2159.
[7] M.T. Darvishi, F. Khani, A series solution of the foam drainage equation, Comput. Math. Appl., 58 (2009) 360-368.
[8] A. Aziz, F. Khani, M.T. Darvishi, Homotopy analysis method for variable thermal conductivity heat flux gage with edge contact resistance, Zeitschrift fuer Naturforschung A, 65a(10) (2010) 771-776.
[9] F. Khani, M.T. Darvishi, R.S.R. Gorla, Analytical investigation for cooling turbine disks with a non-Newtonian viscoelastic fluid, Comput. Math. Appl., 61(7) (2011) 1728-1738.
[10] E. Fan, Z. Jian, Applications of the Jacobi elliptic function method to special-type nonlinear equations, Phys. Lett. A, 305 (6) (2002) 383-392.
[11] M.T. Darvishi, Maliheh Najafi, Mohammad Najafi, Exact three-wave solutions for high nonlinear form of Benjamin-Bona-Mahony-Burgers equations, International Journal of Mathematical and Computer Sciences, 6(3) (2010) 127-131.
[12] M.T. Darvishi, Mohammad Najafi, Some exact solutions of the ( $2+1$ )-dimensional breaking soliton equation using the three-wave method, International Journal of Computational and Mathematical Sciences, 6(1) (2012) 13-16.
[13] M.T. Darvishi, Maliheh Najafi, Mohammad Najafi, New exact solutions for the ( $3+1$ )-dimensional breaking soliton equation, International Journal of Information and Mathematical Sciences, 6(2) (2010) 134-137.
[14] M.T. Darvishi, Maliheh Najafi, Mohammad Najafi, New application of EHTA for the generalized (2+1)-dimensional nonlinear evolution equations, International Journal of Mathematical and Computer Sciences, 6(3) (2010) 132-138.
[15] M.T. Darvishi, Mohammad Najafi, A modification of extended homoclinic test approach to solve the (3+1)-dimensional potential-YTSF equation, Chin. Phys. Lett., 28(4) (2011) 040202.
[16] M.T. Darvishi, Mohammad Najafi, Some complexiton type solutions of the (3+1)-dimensional Jimbo-Miwa equation, International Journal of Computational and Mathematical Sciences, 6(1) (2012) 25-27.
[17] M.T. Darvishi, Maliheh Najafi, Mohammad Najafi, Traveling wave solutions for the (3+1)-dimensional breaking soliton equation by $\left(\frac{G^{\prime}}{G}\right)$-expansion method and modified $F$-expansion method, International Journal of Computational and Mathematical Sciences, 6(2) (2012) 64-69.
[18] J.H. He, M.A. Abdou, New periodic solutions for nonlinear evolution equations using Exp-function method, Chaos, Solitons and Fractals, 34 (2007) 1421-1429.
[19] F. Khani, S. Hamedi-Nezhad, M.T. Darvishi, S.-W. Ryu, New solitary wave and periodic solutions of the foam drainage equation using the Exp-function method, Nonlin. Anal.: Real World Appl., 10 (2009) 1904-1911.
[20] B.-C. Shin, M.T. Darvishi, A. Barati, Some exact and new solutions of the Nizhnik-Novikov-Vesselov equation using the Exp-function method, Comput. Math. Appl., 58(11/12) (2009) 2147-2151.
[21] F. Khani, M.T. Darvishi, A. Farmani, L. Kavitha, New exact solutions of coupled (2+1)-dimensional nonlinear system of Schrödinger equations, ANZIAM Journal, 52 (2010) 110-121.
[22] X.H. Wu, J.H. He, Exp-function method and its application to nonlinear equations, Chaos, Solitons and Fractals, 38(3) (2008) 903-910.
[23] R. Camassa, D. Holm, An integrable shallow water equation with peaked solitons, Phys. Rev. Lett. 71(11) (1993) 1661-1664.
[24] L. Tian, X. Song, New peaked solitary wave solutions of the generalized CamassaHolm equation, Chaos Solitons and Fractals. 19 (2004) 621-637.
[25] J. P. Boyd, Peakons and cashoidal waves: traveling wave solutions of the CamassaHolm equation, Appl. Math. Comput. 81 (23) (1997) 73187.
[26] A. Wazwaz, New compact and noncompact solutions for two variants of a modified Camassa Holm equation, Appl. Math. Comput. 163 (3) (2005) 11651179.
[27] A. Irshad, M. Usman and S. T. Mohyud-Din, Exp-function Method for Simplified Modified Camassa Holm Equation , International Journal of Modern Mathematical Sciences 4(3) (2012) 146-155.
[28] J.H. He, Some asymptotic methods for strongly nonlinear equations, Internat. J. Modern Phys. B, 20 (2006) 1141-1199.
[29] J.H. He, Variational principles for some nonlinear partial differential equations with variable coefficients, Chaos, Solitons and Fractals, 19 (2004) 847-851.

