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On the Role of Strong Coupling Constant and Nucleons in Understanding Nuclear Stability and Binding Energy

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ABSTRACT

In this study, we dealt with beta stability line and nuclear binding energy in a simplified semi empirical approach by taking into account nucleon and electron rest masses and with reference to nuclear charge radius and strong coupling constant. In this context, we have considered $(1/\alpha_s)(e^2/8\pi\epsilon_0R_0) \cong 4.86$ MeV as a single nuclear binding energy coefficient or potential.

Keywords: Nucleon, Nuclear binding energy, Beta stability line, Strong coupling, Semi-empirical mass formula, Liquid drop model

1. Introduction

In nuclear physics, the semi-empirical mass formula [1-3] is used to approximate the mass and various other properties of an atomic nucleus. As its name suggests, it is based partly on theory and partly on empirical measurements. Although refinements have been made to the coefficients over the years, the structure of the formula remains the same today. The simple semi empirical mass formula (SEMF) constitutes of five different energy coefficients and 5 different energy terms. The major drawback of SEMF is that, it does not throw light on the implementation of strong interaction concepts and strong coupling constant [4,5] in understanding nuclear binding energy. In this study, we attempt:

1) To propose a very simple and direct relation for understanding beta stability line [5, 6, 7] with reference to neutron, proton and electron rest masses,

2) To propose a very simple semi empirical relation for fitting the nuclear binding energy with only one energy constant i.e. 4.86 MeV, which is assumed to be

connected with nuclear charge radius $R_0 \approx 1.25$ fm and strong coupling constant $\alpha_s \approx 0.1186$,

3) To divide the nuclear binding energy scheme into 3 steps:

a. To understand the binding energy at stable mass number of Z,

b. To understand the binding energy below the stable mass number of Z,

c. To understand the binding energy above the stable mass number of Z;

4) To report an attempt to fit and interrelate the SEMF binding energy coefficients.

2. Liquid Drop Model and the Semi Empirical Mass Formula

According to the well-known liquid drop model:

1) Atomic nucleus can be considered as a drop of incompressible fluid;

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2) Nuclear fluid is made of protons and neutrons, which are held together by the strong nuclear force.

Mathematical analysis of the theory delivers an equation which attempts to predict the binding energy of a nucleus in terms of the numbers of protons and neutrons it contains. This equation includes five terms on its right hand side. These correspond to the cohesive binding of all the nucleons by the strong nuclear force, the electrostatic mutual repulsion of the protons, a surface energy term, an asymmetry term which is derivable from the protons and neutrons occupying independent quantum momentum states, and a pairing term which is partly derivable from the protons and neutrons occupying independent quantum spin states. The coefficients are calculated by fitting to experimentally measured masses of nuclei. Their values can vary depending on how they are fitted to the data. In the following formulae, let A be the total number of nucleons, Z the number of protons, and N the number of neutrons. The mass of an atomic nucleus is given by:

$$m = Zm_p + Nm_n - \left(B/c^2\right) \tag{1}$$

where m_p and m_n represent the rest mass of a proton and a neutron, respectively, and *B* is the binding energy of the nucleus. The semi-empirical mass formula states that the binding energy will take the following form,

$$B = a_v A - a_s A^{2/3} - a_c \frac{Z(Z-1)}{A^{1/3}} - a_a \frac{(A-2Z)^2}{A} \pm \frac{a_p}{\sqrt{A}}$$
(2)

Here a_v is volume energy coefficient, a_s is the surface energy coefficient, a_c is the coulomb energy coefficient, a_a is the asymmetry energy coefficient and a_p is the pairing energy coefficient. If we consider the sum of the volume energy, surface energy, coulomb energy, asymmetry energy and pairing energy, then the picture of a nucleus as a drop of incompressible liquid roughly accounts for the observed variation of binding energy of the nucleus. As given in Table 1, the currently adopted coefficients in unit of energy (MeV) are used in semi empirical mass formula [3].

 Table 1. Adopted SEMF binding energy coefficients

a_{v}	a_s	a_{c}	a_a	a_p
MeV	MeV	MeV	MeV	MeV
15.78	18.34	0.71	23.21	12.0
15.258	16.26	0.689	22.20	10.08

3. Estimation of Stable Mass Number with Proton Number

By maximizing B(A, Z) with respect to Z, it can be found the number of protons Z of the stable nucleus of atomic weight A as,

$$\begin{cases} Z \approx \frac{A}{2 + (a_c/2a_a)A^{2/3}}.\\ (A - 2Z) \approx \frac{0.4A^2}{A + 200} \end{cases}$$
(3)

This is roughly A/2 for light nuclei, but for heavy nuclei there is an even better agreement with nature. By substituting the above value of Z back into B, one obtains the binding energy as a function of the atomic weight, B(A). The maximization of B(A)/A with respect to A would give the nucleus which is most strongly bound or most stable. In this context, we would like to suggest that, independent of SEMF concepts, nuclear beta stability line can be understood within neutron, proton and electron rest masses. It is also possible to show that [4]:

$$\exp\left(\frac{\left(m_{n}-m_{p}\right)c^{2}}{m_{e}c^{2}}\right) \cong \exp\left(\frac{1.293332\text{MeV}}{0.5109989461\text{MeV}}\right) \cong 4\pi \qquad (4)$$
$$\rightarrow \left(\frac{\left(m_{n}-m_{p}\right)c^{2}}{m_{e}c^{2}}\right) \cong \ln\left(4\pi\right) \cong 2.53102425$$

Based on this observation and without considering the binding energy coefficients, beta stability line can be understood with the following empirical relations.

These relations can be compared with the computational relations pertaining to isotopic shift and drip lines proposed in reference [3], $N_s = 0.968051Z + 0.00658803Z^2$.

Let,
$$k \approx \left(\frac{1}{4\pi}\right)^2 \approx 0.00633$$

 $A_s \approx 2Z + \left(\frac{Z}{4\pi}\right)^2 \approx 2Z + 0.00633Z^2 \approx 2Z + kZ^2$
 $N_s \approx Z + \left(\frac{Z}{4\pi}\right)^2 \approx 2Z + 0.00633Z^2 \approx Z + kZ^2$

$$A_s - 2Z \approx \left(\frac{Z}{4\pi}\right)^2 \approx 0.00633Z^2 \approx kZ^2$$
(5)

With even-odd corrections much better correlations can be observed. For light and medium atomic nuclides, there is some mismatch, which can be attributed to shell structure and needs for further study. As indicated in Table 2 the stable nucleon number can be estimated with its corresponding proton number by fitting.

Proton Number	Estimated Stable Mass	A_s with Even
Ζ	Number, A_s	Odd Correction
21	44.80	45
29	63.32	63
37	82.67	83
47	107.99	107/109
53	123.79	123/125
60	142.80	142
69	168.15	167/169
79	197.52	197
83	209.62	209/211
92	237.60	238
100	263.33	262
112	303.43	302
118	324.17	324

Table 2. The fit results for the stable mass numbersversus proton number

4. Characteristic Nuclear Binding Potential and Nuclear Binding Energy Coefficients

In this analysis, it is assumed that with reference to nuclear charge radius, $R_0 \cong 1.25$ fm, and the strong coupling constant, $\alpha_s \cong 0.1186$ by the following relation.

$$B_0 \cong \left(\frac{1}{\alpha_s}\right) \frac{e^2}{8\pi\varepsilon_0 R_0} \cong 4.85655 \text{ MeV} \cong 4.86 \text{ MeV}$$
(6)

With this binding energy potential, nuclear binding energy can be fitted with the following two-term semi empirical relation. It should be noted that, with further research and analysis, qualitatively a simplified and unified method can be developed. With this energy potential, energy coefficients of the SEMF can be fitted in the following way.

$$\begin{cases} a_{v} \approx \frac{1}{2} \left(\frac{m_{p}}{m_{e}}\right)^{\frac{1}{4}} B_{0} \approx 15.91 \text{ MeV}; \quad a_{s} \approx a_{v} + \frac{1}{2} B_{0} \approx 18.34 \text{ MeV} \\ a_{c} \approx \left(\frac{m_{p}}{m_{e}}\right)^{-\frac{1}{4}} B_{0} \approx 0.7425 \text{ MeV}; \quad a_{a} \approx a_{s} + B_{0} \approx 23.20 \text{ MeV} \end{cases}$$
(7)
$$a_{p} \approx \frac{1}{2} a_{a} \approx \frac{1}{2} (a_{s} + B_{0}) \approx 11.6 \text{ MeV}; \\ (a_{s} - a_{v}) \approx \frac{1}{2} B_{0}; \quad (a_{a} - a_{s}) \approx B_{0}; \quad (a_{a} - a_{v}) \approx \frac{3}{2} B_{0}; \end{cases}$$

5. Proposed Method of Estimating Nuclear Binding Energy for Z=2 to 83

Starting from Z=2 to 83,

Step 1: Close to beta stability line, nuclear binding energy can be expressed with the following semi empirical relation.

$$B_{A_s} \cong x \left[\left(\frac{A_s^2}{A_s - Z} \right) - \frac{\left(A_s - 2Z\right)^2}{A_s} \right] 4.86 \text{ MeV}$$
(8)

where,

$$Z \approx 2$$
 to 29, $x \approx \left(\frac{Z}{30}\right)^{\frac{1}{6}}$ and for $30 \ge Z \le 83, x \approx 1.0$

$$B_A \cong xy \left[\left(\frac{A_s^2}{A_s - Z} \right) - \frac{\left(A_s - 2Z\right)^2}{A_s} \right] 4.86 \text{ MeV} \quad (9)$$

Step 2: Below and above the stable mass numbers,

where,

$$Z \approx 2 \text{ to } 29, \ x \approx \left(\frac{Z}{30}\right)^{\frac{1}{6}}$$

and for $30 \ge Z \le 83, \ x \approx 1.0$
$$\begin{cases} A < A_s, \ y \approx \left(\frac{A - Z}{A_s - Z}\right)^{\frac{3}{2}} \\ A > A_s, \ y \approx \left(\frac{A - Z}{A_s - Z}\right)^{\frac{1}{2}} \end{cases}$$

2)

6. Alternative Expression for Asymmetry Energy Term at Beta Stability Line

With reference to the proposed beta stability line coefficient, $k \cong \left(\frac{1}{4\pi}\right)^2 \cong 0.006333$, close to the beta stability line, with trial error, we noticed that:

For light and medium atoms,

$$\frac{(A_s - 2Z)^2}{A_s} \approx k^2 A_s N_s^{3/2}$$
(10)

For heavy atoms,

$$k^{2}A_{s}N_{s}^{3/2} > \frac{\left(A_{s} - 2Z\right)^{2}}{A_{s}}$$
(11)

7. Proposed Method of Estimating Nuclear Binding Energy for Z > 83

Based on the above observation, for Z > 83, binding energy close to beta stability line can be expressed with the following relation.

$$B_{A_s} \cong x \left[\left(\frac{A_s^2}{A_s - Z} \right) - \left(k^2 A_s N_s^{3/2} \right) \right] 4.86 \text{ MeV} \quad (12)$$

where $x \approx 1.0$

$$B_A \cong xy \left[\left(\frac{A_s^2}{A_s - Z} \right) - \left(k^2 A_s N_s^{3/2} \right) \right] 4.86 \text{ MeV}$$
(13)

Below and above the stable mass number,

where,

$$\begin{cases} x \approx 1.0, \\ A < A_s, \ y \approx \left(\frac{A-Z}{A_s-Z}\right)^{\frac{2}{3}} \text{ and } A > A_s, \ y \approx \left(\frac{A-Z}{A_s-Z}\right)^{\frac{1}{2}} \end{cases}$$

8. Binding Energy Data Fitting

As given in Table 3 the nuclear binding energy was estimated at stable mass numbers.

Fig.1 shows nuclear binding energy at stable atomic nuclides. The plotted solid curve between the binding energy and the proton number indicates the actual or SEMF binding energy (MeV) and dotted red curve indicates estimated binding energy (MeV).



Fig. 1 Nuclear binding energy at stable atomic nuclides

As seen in Table 4, the isotopic binding energies are calculated for Z=20, 30, 40, 50, 60, 70, 80 and 90. The column-1 represents the proton number, column-2 represents the stable estimated mass number, column-3 represents the neutron number, column-4 represents the nucleon number, column-5 represents the binding energy calculated with SEMF, column-6 represents the binding energy calculated with proposed relations and column 7 represents the %error with respect to SEMF.

Table 3. Estimated nuclear binding energy at stable mass numbers A_s from 4 to 263

Proton	Stable Mass	A_{c}	Estimated	Actual [8] or		
Number	Number A_s	with Even	Binding	*SEMF [2]		
Ζ		Odd	Energy (MeV)	Binaing Energy (MeV)		
	1	<u>Correction</u>	24.8	28.206		
2	4	4	24.8	28.290		
3	0	/	40.1	59.244		
4	8	8	55.6	58.08		
5	10	11	72.4	76.205		
6	12	12	89.2	92.162		
7	14	14	106.8	104.659		
8	16	16	124.8	127.619		
9	19	19	143.3	147.801		
10	21	21	162.0	167.406		
11	23	23	181.1	186.564		
12	25	24	200.2	198.257		
13	27	27	220.0	224.952		
14	29	28	239.7	236.537		
15	31	31	259.9	262.917		
16	34	34	280.6	291.839		
17	36	35	300.7	298.21		
18	38	38	321.8	327.343		
19	40	39	342.4	333.724		
20	43	42	363.8	361.896		
21	45	45	385.5	387.848		
22	47	46	406.5	398.193		
23	49	49	428.5	*426.34		
24	52	52	450.8	456.349		
25	54	55	473.2	482.075		
26	56	56	494.7	492.258		
27	59	59	517.5	517.313		
28	61	60	539.2	526.846		
29	63	63	562.2	551.385		
30	66	66	585.4	578.136		
31	68	67	607.5	*582.1		
32	70	70	624.2	610.521		
33	73	73	644.2	*634.34		
34	75	74	663.0	642.891		
35	78	77	683.0	*667.66		
36	80	80	703.0	695.434		
37	83	83	723.1	*720.46		
38	85	84	741.8	728.906		
39	88	87	761.8	*755.05		
40	90	90	781.9	783.893		
41	93	93	802.0	805.765		
42	95	94	820.7	814.256		

43	98	97	840.7	*836.26
44	100	100	860.9	861.928
45	103	103	881.0	884.163
46	105	104	899.6	892.82
47	108	107	919.7	915.263
48	111	110	939.8	940.646
49	113	113	960.0	963.094
50	116	116	980.1	988.684
51	118	117	998.7	*992.78
52	121	120	1018.8	1017.282
53	124	123	1039.0	*1038.77
54	126	126	1059.1	1063 909
55	120	120	1079.3	*1085.08
55	129	132	1099.5	1110.038
57	132	132	11197	*1131
58	135	135	1138.1	1138 702
50	140	130	1158.3	*1160.56
59	140	137	1138.5	1195 142
60	145	142	11/8.3	*1202.96
61	140	145	1198.7	*1203.86
62	148	148	1218.9	1225.392
63	151	151	1239.1	1244.141
64	154	154	1259.3	1266.627
65	157	157	1279.5	*1287.38
66	160	160	1299.8	1309.455
67	162	161	1318.2	*1314.32
68	165	164	1338.4	1336.447
69	168	167	1358.6	*1356.45
70	171	170	1378.8	1378.13
71	174	173	1399.0	*1397.78
72	177	176	1419.3	1418.801
73	180	179	1439.5	*1438.48
74	183	182	1459.7	1459.335
75	186	185	1479.9	1478.341
76	189	186	1498.3	1484.807
77	192	191	1520.4	1518.088
78	195	194	1540.6	1539.577
79	198	197	1560.9	1559.386
80	201	200	1581.1	1581.181
81	204	203	1601.4	1600.87
82	207	206	1621.6	1622.325
83	210	209	1641.8	1640.23
84	213	212	1646.6	*1655.76
85	216	215	1666.0	*1669.2
86	219	218	1685.3	*1685.75
87	2.2.2	221	1704.6	*1701.68
<u> </u>	225	224	1723.8	*1719.13

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89	228	227	1743.0	*1735.53
90	231	230	1762.1	*1753.97
91	234	233	1781.2	*1770.78
92	238	238	1803.7	1801.69
93	241	241	1822.7	*1817.31
94	244	244	1841.6	*1835.45
95	247	247	1860.5	*1851.73
96	250	250	1879.3	*1868.97
97	254	253	1898.0	*1884.38
98	257	256	1916.8	*1901.62
99	260	259	1935.4	*1917.39
100	263	262	1954.0	*1935.12

Table 4. Estimated isotopic nuclear binding energy of Z=20 to 90

Proton Number	Estimated Stable Mass Number	Neutron Number	Nucleon Number	SEMF Binding Energy in MeV	Proposed Binding Energy in MeV	%Difference
20	42	20	40	339.7	341.4	-0.50
20	42	21	41	350.1	352.7	-0.74
20	42	22	42	363.2	363.8	-0.17
20	42	23	43	371.6	372.0	-0.10
20	42	24	44	382.7	380.0	0.71
20	42	25	45	389.3	387.8	0.38
20	42	26	46	398.7	395.5	0.80
20	42	27	47	403.8	403.0	0.19
20	42	28	48	411.8	410.4	0.33
20	42	29	49	415.5	417.7	-0.53
20	42	30	50	422.1	424.8	-0.65
20	42	31	51	424.7	431.8	-1.68
20	42	32	52	430.2	438.8	-1.99
20	42	33	53	431.7	445.6	-3.21
30	66	30	60	509.5	518.4	-1.75
30	66	31	61	521.1	529.9	-1.68
30	66	32	62	535	541.2	-1.16
30	66	33	63	545.2	552.4	-1.32
30	66	34	64	557.8	563.5	-1.02
30	66	35	65	566.7	574.5	-1.38
30	66	36	66	577.9	585.4	-1.30
30	66	37	67	585.7	593.5	-1.33
30	66	38	68	595.8	601.4	-0.95
30	66	39	69	602.4	609.3	-1.15

30	66	40	70	611.5	617.1	-0.91
30	66	41	71	617.2	624.7	-1.22
30	66	42	72	625.3	632.3	-1.12
30	66	43	73	630.1	639.8	-1.54
30	66	44	74	637.2	647.2	-1.57
30	66	45	75	641.2	654.5	-2.07
30	66	46	76	647.6	661.7	-2.18
40	90	40	80	666.2	673.8	-1.14
40	90	41	81	678.6	685.0	-0.94
40	90	42	82	693.1	696.1	-0.43
40	90	43	83	704.3	707.1	-0.40
40	90	44	84	717.7	718.0	-0.05
40	90	45	85	728	728.9	-0.12
40	90	46	86	740.4	739.6	0.11
40	90	47	87	749.7	750.3	-0.08
40	90	48	88	761.2	760.9	0.04
40	90	49	89	769.6	771.4	-0.24
40	90	50	90	780.2	781.9	-0.22
40	90	51	91	787.8	789.7	-0.24
40	90	52	92	797.6	797.4	0.03
40	90	53	93	804.4	805.0	-0.08
40	90	54	94	813.4	812.6	0.10
40	90	55	95	819.5	820.1	-0.07
40	90	56	96	827.8	827.5	0.04
40	90	57	97	833.3	834.8	-0.18
40	90	58	98	840.8	842.1	-0.16
40	90	59	99	845.7	849.4	-0.43
50	116	56	106	887.6	878.4	1.03
50	116	57	107	898.1	888.8	1.03
50	116	58	108	910.4	899.2	1.23
50	116	59	109	920.1	909.5	1.15
50	116	60	110	931.8	919.8	1.29
50	116	61	111	940.7	930.0	1.14
50	116	62	112	951.6	940.1	1.21
50	116	63	113	960.0	950.2	1.02
50	116	64	114	970.2	960.2	1.03
50	116	65	115	977.9	970.2	0.79
50	116	66	116	987.5	980.1	0.75
50	116	67	117	994.6	987.5	0.71
50	116	68	118	1003.6	994.8	0.87
50	116	69	119	1010.1	1002.1	0.79
50	116	70	120	1018.5	1009.4	0.90
50	116	71	121	1024.4	1016.5	0.77

50	116	72	122	1032.3	1023.7	0.83
50	116	73	123	1037.8	1030.8	0.68
50	116	74	124	1045.1	1037.8	0.70
50	116	75	125	1050.1	1044.8	0.51
50	116	76	126	1056.9	1051.7	0.49
60	142	73	133	1099	1090.6	0.76
60	142	74	134	1110.2	1100.5	0.87
60	142	75	135	1119	1110.4	0.77
60	142	76	136	1129.7	1120.3	0.83
60	142	77	137	1138	1130.1	0.69
60	142	78	138	1148	1139.9	0.71
60	142	79	139	1155.8	1149.6	0.54
60	142	80	140	1165.4	1159.3	0.53
60	142	81	141	1172.7	1168.9	0.32
60	142	82	142	1181.7	1178.5	0.27
60	142	83	143	1188.5	1185.7	0.24
60	142	84	144	1197.1	1192.8	0.36
60	142	85	145	1203.5	1199.9	0.30
60	142	86	146	1211.6	1206.9	0.39
60	142	87	147	1217.5	1213.9	0.30
60	142	88	148	1225.2	1220.9	0.35
60	142	89	149	1230.7	1227.8	0.24
60	142	90	150	1237.9	1234.7	0.26
60	142	91	151	1243	1241.5	0.12
60	142	92	152	1249.9	1248.3	0.13
60	142	93	153	1254.6	1255.1	-0.04
70	170	91	161	1303.9	1294.8	0.70
70	170	92	162	1313.9	1304.2	0.73
70	170	93	163	1321.8	1313.7	0.61
70	170	94	164	1331.3	1323.1	0.62
70	170	95	165	1338.8	1332.4	0.47
70	170	96	166	1348.0	1341.8	0.46
70	170	97	167	1355.0	1351.1	0.29
70	170	98	168	1363.8	1360.4	0.25
70	170	99	169	1370.5	1369.6	0.07
70	170	100	170	1378.8	1378.8	0.00
70	170	101	171	1385.1	1385.7	-0.04
70	170	102	172	1393.0	1392.5	0.03
70	170	103	173	1399.0	1399.3	-0.02
70	170	104	174	1406.6	1406.1	0.04
70	170	105	175	1412.1	1412.8	-0.05
70	170	106	176	1419.4	1419.6	-0.01
70	170	107	177	1424.6	1426.2	-0.12

70	170	108	178	1431.5	1432.9	-0.10
70	170	109	179	1436.4	1439.5	-0.22
70	170	110	180	1443.0	1446.1	-0.21
70	170	111	181	1447.5	1452.7	-0.36
00	200	111	101	1500.0	15010	0.50
80	200	111	191	1509.8	1501.0	0.58
80	200	112	192	1518.7	1510.0	0.57
80	200	113	193	1525.6	1519.0	0.43
80	200	114	194	1534.1	1527.9	0.40
80	200	115	195	1540.8	1536.9	0.26
80	200	110	196	1549.0	1545.8	0.21
80	200	11/	197	1555.5	1554.0	0.04
80	200	118	198	1503.1	1503.5	-0.02
80 80	200	119	199 200	1509.1	1572.5	-0.20
ðU 80	200	120	200	15/0./	1581.1	-0.28
80	200	121	201	1580.6	1504.2	-0.34
80	200	122	202	1505.0	1594.2	-0.29
80	200	125	203	1595.0	1607.2	-0.30
80	200	124	204	1607.0	1612.7	-0.33
80	200	125	205	1612.6	1620.1	-0.42
80	200	120	200	1618.5	1620.1	-0.41
80	200	127	207	1618.3	1620.0	-0.30
80	200	120	208	1620.4	1633.0	-0.30
80	200	129	209	1625.4	1645.7	-0.61
80	200	121	210	1630.7	1652.0	-0.03
80	200	151	211	1039.7	1052.0	-0.75
90	230	131	221	1699.8	1685.7	0.83
90	230	132	222	1707.9	1694.3	0.80
90	230	133	223	1714.2	1702.9	0.66
90	230	134	224	1722.1	1711.4	0.62
90	230	135	225	1728.1	1719.9	0.48
90	230	136	226	1735.6	1728.4	0.42
90	230	137	227	1741.4	1736.8	0.26
90	230	138	228	1748.7	1745.3	0.20
90	230	139	229	1754.2	1753.7	0.03
90	230	140	230	1761.2	1762.1	-0.05
90	230	141	231	1766.5	1768.4	-0.11
90	230	142	232	1773.2	1774.6	-0.08
90	230	143	233	1778.3	1780.9	-0.15
90	230	144	234	1784.8	1787.1	-0.13
90	230	145	235	1789.5	1793.3	-0.21
90	230	146	236	1795.8	1799.5	-0.20
90	230	147	237	1800.3	1805.6	-0.30
90	230	148	238	1806.3	1811.7	-0.30

90	230	149	239	1810.6	1817.9	-0.40
90	230	150	240	1816.4	1823.9	-0.42
90	230	151	241	1820.5	1830.0	-0.52

9. Discussion

From the above proposed relations and the estimated results in tables 2 to 4 it is possible to infer that:

1) Nucleons and electrons play a crucial role in understanding nuclear stability.

2) Strong coupling constant plays an important role in understanding the nuclear binding energy.

3) Nuclear binding energy can also be estimated with use of a single energy coefficient.

4) Z=30 seems to play an interesting role and we are working on understanding the physical significance of

$$\left(\frac{Z}{30}\right)^{\frac{1}{6}}$$
.

5) Z=53 is estimated to be stable $A_s=123$ and its estimated binding energy is 1039 MeV. Actually, it is stable at $A_s=127$. From relation (9), binding energy of

 $_{53}I^{127}$ is $\left(\frac{127-53}{123-53}\right)^{\frac{1}{2}}$ 1039.0 \cong 1068.27 MeV. Actual binding energy of $_{53}I^{127}$ is 1072.57 MeV.

6) The authors are working in this line to deal with the nuclear binding energies and masses involving more realistic quantum physics, different shell corrections and odd-even phenomena [9-14].

10. Conclusion

Understanding and estimating nuclear binding energy with strong interaction concepts still stands as a really challenging task. So far no such model is available in physics literature. Even though some % difference is persisting in the proposed binding energy expressions, qualitatively they are very simple to follow. An interesting point to be noted is that, the number of energy coefficients can be minimized and their existence can be linked with strong interaction concepts. We are confident to say that, with further research, background physics of the proposed relations can be understood and thereby a clear and simple model can be developed.

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Conflicts of Interest

The authors have no conflict of interest.

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