



A new formulation for the linearized Navier-Stokes equation

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Abstract

This paper is devoted to study the Navier-Stokes equations by applying the curl and using a current function, we obtain a non-linear biharmonic problem where the pressure disappears and instead of the velocity, we are working with a scalar function. After a linearization, we obtain a sequence of linear problems. We study the existence and uniqueness of its solutions. Finally we show the convergence of the sequence of the linearized problems obtained to the non-linear one.

Keywords: Bi-Laplacian, Existence and uniqueness, Navier-Stokes equations.

1. Introduction

We consider the Navier-Stokes problem:

$$(P) \begin{cases} -\nu \Delta u + (u \nabla) u + \nabla P = f & \text{in } \Omega, \\ \operatorname{div} u = 0 & \text{in } \Omega, \\ u = 0 & \text{on } \Gamma, \end{cases}$$

where Ω is a bounded and connected domain in \mathbb{R}^2 with lipschitz boundary $\Gamma = \partial\Omega$,

\mathbf{u} the velocity and p the pressure.

ν is a positive parameter called kinematic viscosity and the corresponding function \mathbf{f} forces applied to the fluid is given.

After the application of the curl and using a current function, we obtain a non-linear biharmonic problem.

The variational formulation of the Navier-Stokes equations in the classic form is well studied in [2], [4], and [6]. A discretisation by Finite Element Methods of the problem is proposed by [4].

The standard discretization of the Stokes and Navier-Stokes equations in vorticity and stream function formulation by affine finite elements is known for its bad convergence. Amara.M and Bernardi.C in [1] present a modified discretization, they prove that the convergence is improved and they establish a priori error estimates.

The outline of the paper is as follows:

In the second Section, we are concerned with the bi-harmonic equation by applying the curl and using a current function.

In Section 3, we study the sequence of linearized problems. We show the existence and uniqueness of their solutions .

In Section 4, we show the convergence of the sequence of solutions of the linearized problems obtained to the non-linear one.

In section 5, we demonstrated the linear convergence.

2. Application of rotational

We have $\operatorname{div} u = 0$, then it can be written in the form $u = \operatorname{curl} \phi$ where ϕ is a scalar function called fairly regular stream function.

$$\begin{cases} u_1 = \frac{\partial \phi}{\partial y}, \\ u_2 = -\frac{\partial \phi}{\partial x}. \end{cases}$$

Then by applying the curl to our problem (P), we will have:

$$\operatorname{curl} \Delta u = -\Delta^2 \phi,$$

$$\operatorname{curl} (\nabla P) = 0,$$

$$\operatorname{curl}((u \cdot \nabla)u) = \frac{\partial}{\partial x}(u_1 \frac{\partial u_2}{\partial x} + u_2 \frac{\partial u_2}{\partial y}) - \frac{\partial}{\partial y}(u_1 \frac{\partial u_1}{\partial x} + u_2 \frac{\partial u_1}{\partial y}),$$

since

$$\operatorname{div} u = 0 \quad \text{this implies that} \quad \frac{\partial u_1}{\partial x} = -\frac{\partial u_2}{\partial y},$$

and then, we obtain:

$$\operatorname{curl}((u \cdot \nabla)u) = -\frac{\partial \phi}{\partial y} \frac{\partial \Delta \phi}{\partial x} + \frac{\partial \phi}{\partial x} \frac{\partial \Delta \phi}{\partial y}.$$

The equation becomes

$$\nu \Delta^2 \phi - \frac{\partial \phi}{\partial y} \frac{\partial \Delta \phi}{\partial x} + \frac{\partial \phi}{\partial x} \frac{\partial \Delta \phi}{\partial y} = \operatorname{curl} f,$$

and we have the following problem

$$(Q) \begin{cases} \nu \Delta^2 \phi - \frac{\partial \phi}{\partial y} \frac{\partial \Delta \phi}{\partial x} + \frac{\partial \phi}{\partial x} \frac{\partial \Delta \phi}{\partial y} = \operatorname{curl} f & \text{in } \Omega, \\ \phi = 0 & \text{on } \Gamma, \\ \frac{\partial \phi}{\partial x} = \frac{\partial \phi}{\partial y} = 0 & \text{on } \Gamma. \end{cases}$$

A linearization gives:

$$\nu \Delta^2 \phi_{n+1} - \frac{\partial \phi_n}{\partial y} \frac{\partial \Delta \phi_{n+1}}{\partial x} + \frac{\partial \phi_n}{\partial x} \frac{\partial \Delta \phi_{n+1}}{\partial y} = \operatorname{curl} f.$$

We set

$$c_n = \frac{\partial \phi_n}{\partial y}, \quad d_n = \frac{\partial \phi_n}{\partial x},$$

we have:

$$\nu \Delta^2 \phi_{n+1} - c_n \frac{\partial \Delta \phi_{n+1}}{\partial x} + d_n \frac{\partial \Delta \phi_{n+1}}{\partial y} = \text{curl } f.$$

Therefore, our problem is:

$$(Q_{n+1}) \begin{cases} \nu \Delta^2 \phi_{n+1} - c_n \frac{\partial \Delta \phi_{n+1}}{\partial x} + d_n \frac{\partial \Delta \phi_{n+1}}{\partial y} = \text{curl } f & \text{in } \Omega, \\ \phi_{n+1} = 0 & \text{on } \Gamma, \\ \frac{\partial \phi_{n+1}}{\partial x} = \frac{\partial \phi_{n+1}}{\partial y} = 0 & \text{on } \Gamma, \end{cases}$$

3. Variational formulation

We multiply both sides of the first equation of (Q_{n+1}) by a test function $v \in V = H_0^2(\Omega)$ and integrating over Ω , we have the following variational problem:

$$(QV)_{n+1} \begin{cases} \text{Find } \phi_{n+1} \in V \text{ such as} \\ a(\phi_{n+1}, v) = L(v), \quad \forall v \in V, \end{cases}$$

where $a(\cdot, \cdot)$ is a bilinear form on $V \times V$ given by:

$$a(u, v) = a_0(u, v) + a_n(u, v), \tag{1}$$

where

$$a_0(u, v) = \nu \int_{\Omega} \Delta u \Delta v dX, \tag{2}$$

$$a_n(u, v) = \int_{\Omega} \Delta u \left(\frac{\partial \phi_n}{\partial y} \frac{\partial v}{\partial x} - \frac{\partial \phi_n}{\partial x} \frac{\partial v}{\partial y} \right) dX, \tag{3}$$

and $L(\cdot)$ is a linear form on V defined by the following expression:

$$L(v) = \int_{\Omega} \text{curl } f v dX. \tag{4}$$

For the existence and uniqueness of the solution, we need this lemma:

Lemma 3.1. [5], $\forall u \in H^2(\Omega) \cap H_0^1(\Omega)$, there exist $c'(\Omega) > 0$ such that

$$\|u\|_{H^2}^2 \leq c'(\Omega) \|\Delta u\|_2^2 \leq c^* \|\Delta u\|_2^2. \tag{5}$$

where c^* will be chosen later.

Theorem 3.2. For $f \in H^1(\Omega)$ and $\text{curl } f$ small enough: $\|\text{curl } f\|_2 < c(\nu)$, the problem $(QV)_{n+1}$ has a unique solution $\phi_{n+1} \in V$.

Proof. 1. **Continuity of a :**

We show the continuity and coercivity of the bilinear form $a(\cdot, \cdot)$.

We put $u = \phi_{n+1}$ and we deal with each term separately, we begin by:

$$\begin{aligned} |a_0(u, v)| &\leq \nu \|\Delta u\|_2 \|\Delta v\|_2, \\ &\leq \nu \|u\|_{H^2} \|v\|_{H^2}. \end{aligned} \tag{6}$$

And using the continuous injection of $H^1(\Omega)$ in $L^4(\Omega)$, $a_n(\cdot, \cdot)$ will be bounded as bellow

$$\begin{aligned} |a_n(u, v)| &= \left| \int_{\Omega} \Delta u (\nabla v \wedge \nabla \phi_n) dX \right|, \\ &\leq \|\Delta u\|_2 \|\nabla v\|_{L^4} \|\nabla \phi_n\|_{L^4}, \\ &\leq c \|\Delta u\|_2 \|\nabla v\|_{H^1} \|\nabla \phi_n\|_{H^1}, \\ &\leq c \|\phi_n\|_{H^2} \|u\|_{H^2} \|v\|_{H^2}. \end{aligned} \tag{7}$$

Therefore,

$$|a(u, v)| \leq C_n \|u\|_{H^2} \|v\|_{H^2}, \tag{8}$$

where C_n is a constant which depends on n given by:

$$C_n = \nu + c\|\phi_n\|_{H^2}.$$

This implies that, for each fixed n , $a(., .)$ is continuous on V .

2. Coercivity of a :

We have:

$$a_0(u, u) = \nu \|\Delta u\|_2^2, \tag{9}$$

and taking $v = u$ in (7), we obtain:

$$|a_n(u, u)| \leq c\|\phi_n\|_{H^2} \|u\|_{H^2}^2, \tag{10}$$

then using (5)

$$\begin{aligned} a(u, u) &\geq \nu \|\Delta u\|_2^2 - c\|u\|_{H^2}^2 \|\phi_n\|_{H^2}, \\ &\geq \frac{\nu}{c'(\Omega)} \|u\|_{H^2}^2 - c\|u\|_{H^2}^2 \|\phi_n\|_{H^2}, \\ &\geq \left(\frac{\nu}{c'(\Omega)} - c\|\phi_n\|_{H^2}\right) \|u\|_{H^2}^2. \end{aligned} \tag{11}$$

To get the coercivity, we should have:

$$\frac{\nu}{c'(\Omega)} - c\|\phi_n\|_{H^2} > 0, \quad \forall n \in \mathbb{N}, \tag{12}$$

which means that

$$\|\phi_n\|_{H^2} < \frac{\nu}{c'(\Omega)c}, \quad \forall n \in \mathbb{N}. \tag{13}$$

Let

$$\alpha^* = \frac{\nu}{2c'(\Omega)c}. \tag{14}$$

We take $\phi_0 \in B_{\alpha^*}$ where $B_\alpha = \{v \in H_0^2(\Omega); \|v\|_{H^2} \leq \alpha\}$ and assume that $\phi_n \in B_{\alpha^*}$, we must show by induction that: $\phi_{n+1} \in B_{\alpha^*}$.

Indeed, we have, if we put $u = \phi_{n+1}$:

$$a(u, u) = L(u) = \int_{\Omega} \text{curl } f \cdot u \, dX,$$

$$a(u, u) \leq \|\text{curl } f\|_2 \|u\|_2, \tag{15}$$

then

$$\left(\frac{\nu}{c'(\Omega)} - c\|\phi_n\|_{H^2}\right) \|u\|_{H^2}^2 \leq \|\text{curl } f\|_2 \|u\|_{H^2},$$

and

$$\left(\frac{\nu}{c'(\Omega)} - c\|\phi_n\|_{H^2}\right) \|u\|_{H^2} \leq \|\text{curl } f\|_{L^2}, \tag{16}$$

$$\|u\|_{H^2} \leq \frac{\|\operatorname{curl} f\|_2}{\frac{\nu}{c'(\Omega)} - c\|\phi_n\|_{H^2}}.$$

If we assume that:

$$\|\operatorname{curl} f\|_2 < \frac{\nu^2}{4c'(\Omega)^2c},$$

then we have :

$$\|u\|_{H^2} \leq \frac{\|\operatorname{curl} f\|_2}{\frac{\nu}{c'(\Omega)} - c\alpha^*} \leq \alpha^*,$$

which implies that $a(.,.)$ is coercive on V .

3. Continuity of L :

In the other hand, the linear form L is continuous:

$$\|L(v)\|_2 \leq \|\operatorname{curl} f\|_2 \|v\|_2. \quad (17)$$

Then, using Lax-Milgram Theorem, the problem $(QV)_{n+1}$ has a unique solution $\phi_{n+1} \in V$. □

4. Convergence of the sequence

The sequence $(\phi_n)_{n \in \mathbb{N}}$ obtained in the preceding section verifies:

$$\|\phi_n\|_V \leq \alpha^*, \quad \forall n \geq 0,$$

which implies that the sequence $(\phi_n)_{n \in \mathbb{N}}$ is bounded in V .

Then there exist a subsequence that converges weakly to w in V .

Since the injection of V in $H_0^1(\Omega)$ is continuous, there exists a subsequence still noted ϕ_n which converges strongly to w in $H_0^1(\Omega)$.

For the convergence, we need this regularity result:

Lemma 4.1. *Assume that Ω is of class \mathcal{C}^2 and $\phi_0 \in H_0^2(\Omega) \cap H^3(\Omega)$, we have:*

$$\forall n \in \mathbb{N}, \quad \phi_{n+1} \in H_0^2(\Omega) \cap H^3(\Omega).$$

Proof. The problem (Q_{n+1}) can be written as

$$\begin{cases} -\nu \Delta \omega_{n+1} + c_n \frac{\partial \omega_{n+1}}{\partial x} - d_n \frac{\partial \omega_{n+1}}{\partial y} = -\operatorname{curl} f & \text{in } \Omega, & (1) \\ \Delta \phi_{n+1} = \omega_{n+1} & \text{in } \Omega, & (2) \\ \phi_{n+1} \in H_0^2(\Omega). \end{cases}$$

The variational formulation of (1) is

$$\begin{cases} \text{Find } \omega_{n+1} \in H_0^1(\Omega) \text{ such as} \\ A_n(\omega_{n+1}, v) = l(v), \quad \forall v \in H_0^1(\Omega), \end{cases}$$

where

$$A_n(\omega_{n+1}, v) = \nu \int_{\Omega} \nabla \omega_{n+1} \nabla v \, dX + \int_{\Omega} \omega_{n+1} \left(d_n \frac{\partial v}{\partial y} - c_n \frac{\partial v}{\partial x} \right) dX,$$

$$l(v) = - \int_{\Omega} \operatorname{curl} f \, v \, dX.$$

For the coercivity of A_n , we have:

$$\begin{aligned} A_n(v, v) &= \nu \int_{\Omega} |\nabla v|^2 dX + \int_{\Omega} v (\nabla v \wedge \nabla \omega_n) dX \\ &\geq \nu \|\nabla v\|_2^2 - c \|\nabla v\|_2 \|v\|_{H^1} \|\nabla \omega_n\|_{H^1} \\ &\geq \nu c_1 \|v\|_{H^1}^2 - c\alpha^* \|v\|_{H^1}^2. \end{aligned}$$

It suffices to choose in inequality (5): $c^* > \frac{1}{2c_1}$.

And by Lax-Milgram Theorem, we have $\omega_{n+1} \in H^1(\Omega)$.

The theory of regularity for weak solutions of the laplace problem applied to the variational formulation of (2) gives $\phi_{n+1} \in H^3(\Omega)$. □

Then, we have this result of convergence :

Lemma 4.2. 1. We have, $\forall v \in H_0^2(\Omega) \cap H^3(\Omega)$:

$$\lim_{n \rightarrow +\infty} a_n(\phi_{n+1}, v) = a_0(w, v).$$

2. We have, $\forall v \in H_0^2(\Omega) \cap H^3(\Omega)$:

$$\lim_{n \rightarrow +\infty} a_n(\phi_{n+1}, v) = a_{\infty}(w, v) = \int_{\Omega} \Delta w (\nabla v \wedge \nabla w) dX.$$

Proof. 1. We have:

$$\begin{aligned} |a_0(\phi_{n+1}, v) - a_0(w, v)| &= \left| \nu \int_{\Omega} \Delta(\phi_{n+1} - w) \Delta v dX \right| \\ &\leq \nu \|\nabla(\phi_{n+1} - w)\|_2 \|v\|_{H^3}, \end{aligned}$$

then, we obtain the result.

2. On the other hand, we have:

$$|a_n(\phi_{n+1}, v) - a_{\infty}(w, v)| = T_{1_n} + T_{2_n}$$

where

$$T_{1_n} = \left| \int_{\Omega} (\Delta\phi_{n+1} - \Delta w)(\nabla v \wedge \nabla\phi_n) dX \right|,$$

and

$$T_{2_n} = \left| \int_{\Omega} \Delta w \cdot [\nabla v \wedge (\nabla\phi_n - \nabla w)] dX \right|$$

By Green formula we have :

$$T_{1_n} = \left| \int_{\Omega} (\nabla\phi_{n+1} - \nabla w) \nabla(\nabla v \wedge \nabla\phi_n) dX \right|$$

which gives :

$$T_{1_n} \leq \sum_{i,j=1}^2 \left| \int_{\Omega} \frac{\partial(\phi_{n+1} - \omega)}{\partial x_i} \frac{\partial^2 v}{\partial x_i \partial x_j} \frac{\partial(\phi_n)}{\partial x_j} dX \right| + \sum_{i,j=1}^2 \left| \int_{\Omega} \frac{\partial(\phi_{n+1} - \omega)}{\partial x_i} \frac{\partial v}{\partial x_j} \frac{\partial^2(\phi_n)}{\partial x_i \partial x_j} dX \right|.$$

and

$$T_{2_n} \leq \sum_{i,j=1}^2 \left| \int_{\Omega} \frac{\partial(\phi_n - \omega)}{\partial x_i} \frac{\partial^2 \omega}{\partial x_i \partial x_j} \frac{\partial v}{\partial x_j} dX \right|$$

According to the Sobolev imbedding Theorem the space $H^1(\Omega)$ is continuously imbedded in $L^4(\Omega)$ for $n = 2$.

Then by the Hölder's inequality we have for $\omega, v, \phi_n \in H_0^2(\Omega) \cap H^3(\Omega)$:

$$\frac{\partial \omega}{\partial x_i} \frac{\partial^2 v}{\partial x_i \partial x_j} \frac{\partial(\phi_{n+1} - \omega)}{\partial x_j} \in L^1(\Omega), \quad 1 \leq i, j \leq 2,$$

with for T_{1n}

$$\begin{aligned} \left| \int_{\Omega} \frac{\partial \phi_n}{\partial x_i} \frac{\partial^2 v}{\partial x_i \partial x_j} \frac{\partial(\phi_{n+1} - \omega)}{\partial x_j} dX \right| &\leq \left\| \frac{\partial^2 v}{\partial x_i \partial x_j} \right\|_4 \left\| \frac{\partial \phi_n}{\partial x_i} \right\|_4 \left\| \frac{\partial(\phi_{n+1} - \omega)}{\partial x_j} \right\|_2 \\ &\leq C \left\| \frac{\partial^2 v}{\partial x_i \partial x_j} \right\|_{H^1} \left\| \frac{\partial \phi_n}{\partial x_i} \right\|_{H^1} \left\| \frac{\partial(\phi_{n+1} - \omega)}{\partial x_j} \right\|_2 \\ &\leq C \|v\|_{H^2} \|\phi_n\|_{H^2} \left\| \frac{\partial(\phi_n - \omega)}{\partial x_j} \right\|_2, \end{aligned}$$

and

$$\begin{aligned} \left| \int_{\Omega} \frac{\partial^2 \phi_n}{\partial x_i \partial x_j} \frac{\partial v}{\partial x_j} \frac{\partial(\phi_{n+1} - \omega)}{\partial x_i} dX \right| &\leq \left\| \frac{\partial^2 \phi_n}{\partial x_i \partial x_j} \right\|_4 \left\| \frac{\partial v}{\partial x_j} \right\|_4 \left\| \frac{\partial(\phi_{n+1} - \omega)}{\partial x_i} \right\|_2 \\ &\leq C \left\| \frac{\partial^2 \phi_n}{\partial x_i \partial x_j} \right\|_{H^1} \left\| \frac{\partial v}{\partial x_j} \right\|_{H^1} \left\| \frac{\partial(\phi_{n+1} - \omega)}{\partial x_i} \right\|_2 \\ &\leq C \|v\|_{H^2} \|\phi_n\|_{H^2} \left\| \frac{\partial(\phi_{n+1} - \omega)}{\partial x_i} \right\|_2, \end{aligned}$$

and for T_{2n} we have

$$\begin{aligned} \left| \int_{\Omega} \frac{\partial^2 \omega}{\partial x_i \partial x_j} \frac{\partial v}{\partial x_j} \frac{\partial(\phi_n - \omega)}{\partial x_i} dX \right| &\leq \left\| \frac{\partial^2 \omega}{\partial x_i \partial x_j} \right\|_4 \left\| \frac{\partial v}{\partial x_j} \right\|_4 \left\| \frac{\partial(\phi_n - \omega)}{\partial x_i} \right\|_2 \\ &\leq C \left\| \frac{\partial^2 \omega}{\partial x_i \partial x_j} \right\|_{H^1} \left\| \frac{\partial v}{\partial x_j} \right\|_{H^1} \left\| \frac{\partial(\phi_n - \omega)}{\partial x_i} \right\|_2 \\ &\leq C \|v\|_{H^2} \|\omega\|_{H^2} \left\| \frac{\partial(\phi_n - \omega)}{\partial x_i} \right\|_2, \end{aligned}$$

Then, by the strongly convergence in $H_0^1(\Omega)$, we will have

$$\lim_{n \rightarrow +\infty} a_n(\phi_{n+1}, v) = a_{\infty}(w, v).$$

□

Proposition 1. w is a solution of Q .

Proof. It follows from Lemma (4.1) that:

$$\lim_{n \rightarrow +\infty} a_0(\phi_{n+1}, v) + a_n(\phi_{n+1}, v) = a_0(w, v) + a_{\infty}(w, v) = L(v). \quad (18)$$

Which gives:

$$\nu \int_{\Omega} \Delta w \Delta v dX + \int_{\Omega} \Delta w (\nabla v \wedge \nabla w) dX = \int_{\Omega} \text{curl } f v dX, \quad (19)$$

then

$$\int_{\Omega} (\nu \Delta^2 w - \frac{\partial w}{\partial y} \frac{\partial \Delta w}{\partial x} + \frac{\partial w}{\partial x} \frac{\partial \Delta w}{\partial y} - \text{curl } f) v dX = 0, \quad \forall v \in H_0^2(\Omega), \quad (20)$$

then

$$\nu \Delta^2 w - \frac{\partial w}{\partial y} \frac{\partial \Delta w}{\partial x} + \frac{\partial w}{\partial x} \frac{\partial \Delta w}{\partial y} = \text{curl } f, \quad \forall w \in H_0^2(\Omega). \quad (21)$$

So we can conclude that w is the solution of Q .

□

5. Linear convergence

For this part, we set

$$w_{n+1} = \phi_{n+1} - w,$$

we have the following result:

Proposition 2. For $w_{n+1} \in H^2(\Omega)$, there exist a constant C such that:

$$\|w_{n+1}\|_{H^2} \leq C\|w_n\|_{H^2}. \quad (22)$$

Proof. Taking the difference between the problem (Q_{n+1}) and the problem (Q) , we have:

$$\nu \Delta^2 w_{n+1} - \frac{\partial \phi_n}{\partial y} \frac{\partial(\Delta w_{n+1})}{\partial x} + \frac{\partial \phi_n}{\partial x} \frac{\partial(\Delta w_{n+1})}{\partial y} = F_n, \quad (23)$$

where

$$F_n = \frac{\partial(\Delta w)}{\partial x} \frac{\partial w_n}{\partial y} - \frac{\partial(\Delta w)}{\partial y} \frac{\partial w_n}{\partial x}. \quad (24)$$

The variational formulation gives us:

$$\begin{aligned} \beta \|w_{n+1}\|_{H^2}^2 &\leq a(w_{n+1}, w_{n+1}) \leq \int_{\Omega} |F_n| w_{n+1} dX, \\ &= \int_{\Omega} |\Delta w (\nabla w_n \wedge \nabla w_{n+1})| dX, \\ &\leq \|\Delta w\|_2 \|\nabla w_n\|_4 \|\nabla w_{n+1}\|_4, \\ &\leq c \|\Delta w\|_2 \|\nabla w_n\|_{H^1} \|\nabla w_{n+1}\|_{H^1}, \\ &\leq c' \|w_n\|_{H^2} \|w_{n+1}\|_{H^2}, \end{aligned}$$

then, we obtain:

$$\|w_{n+1}\|_{H^2} \leq C \|w_n\|_{H^2}, \quad (25)$$

where C is given by:

$$C = \frac{c'}{\beta}, \quad (26)$$

and β the constant of the coercivity.

Which implies the linear convergence. \square

6. Conclusion

In this paper, we studied the Navier-Stokes equations by applying the curl and using a current function through the application of rotational, we obtained a non-linear biharmonic problem.

After a linearization, we proved the existence and uniqueness of weak solution of the variational formulation using Lax-Milgram Theorem and which we can compute by finite element method.

And in a second part we showed the convergence of the sequence as well as the linear convergence.

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