

# A numerical method based on explicit finite difference for solving fractional hyperbolic PDE's

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## Abstract

In this paper, a new numerical scheme based on explicit finite difference approximation for solving fractional hyperbolic partial differential equations (FHPDE's) is formulated. Numerical studies for the model problems are presented to confirm the accuracy and the effectiveness of the proposed method. The obtained results of proposed system are compared with exact solutions and the original system to show the efficient of the new method.

**Keywords:** Fractional Hyperbolic Partial Differential Equations; Preconditioned Explicit Finite Difference Method.

## 1. Introduction

Fractional Partial Differential Equations are extensively used in engineering, physics and mathematical fields, such as porous media, anomalous diffusion, Hamiltonian chaos systems, bioengineering ([1], [2], [3], [4]). They are involved to model physical processes referring to memory properties, genetic characters and path dependence. The fractional hyperbolic partial differential equations (FHPDE's) model the vibrations of structures (e.g. buildings, beams and machines) and are the basis for fundamental equations of atomic physics ([5], [6], [7]). However, only a very few fractional differential equations can be solved analytically because of their complicated form. Hence, the recent rapid development of numerical methods for fractional differential equations has attracted more and more attentions from researchers. The Iterative methods based on the finite difference approximations have been shown suitable for solving the partial differential equations ([8], [9], [10], [11]). Meerschaert and Tadejeran ([12], [13]) proposed a finite difference approximation of fractional advection-diffusion flow equation and two-sided space-fractional differential equation. In this work, the new preconditioned explicit finite difference approximation will be formulated and applied for solving FHPDE's. The structure of this paper is as follows: Section 2 describes the formulation of the preconditioned explicit finite difference method iterative method for solving the fractional hyperbolic partial differential equations. In section 3, the stability analysis of the proposed method will be discussed. Numerical results are presented in order to show the efficiency of the proposed method in section 4. Finally, relevant conclusions are drawn in section 5.

## 2. Formulation of the preconditioned explicit finite difference method for solving the fhpdе's

Consider the fractional order partial differential equations

$$\frac{\partial^2 u(x,t)}{\partial t^2} = c(x,t) \frac{\partial^\alpha u(x,t)}{\partial x^\alpha} + s(x,t), \quad L \leq x \leq R, \quad 0 \leq t \leq T \quad (1)$$

Together with the initial and zero Dirichlet boundary conditions:

$$\left. \begin{aligned} u(x,0) &= f(x), \quad u_t(x,0) = h(x), \quad L \leq x \leq R \\ u(L,t) &= 0, \quad u(R,t) = 0, \quad 0 \leq t \leq T \end{aligned} \right\} \quad (2)$$

Where  $\frac{\partial^\alpha u(x,t)}{\partial x^\alpha}$  denote the left-hand partial fractional derivative of order  $\alpha$  of the function  $u$  with respect to  $x$  and  $1 \leq \alpha \leq 2$ . Nishimoto [14] estimated the left-handed shifted and the right-handed shifted to the left-handed and right-handed derivatives as the following:

$$\frac{d^+ f(x)}{d_+ x^\alpha} = \frac{1}{(\Delta x)^\alpha} \sum_{k=0}^n g_k f(x - (k-1)\Delta x),$$

$$\frac{d^- f(x)}{d_- x^\alpha} = \frac{1}{(\Delta x)^\alpha} \sum_{k=0}^n g_k f(x + (k-1)\Delta x)$$

Where  $n$  is the number of subdivision of interval  $[L, R]$  and  $\alpha$  is the fractional number. Therefore, we can write:

$$\begin{aligned} \frac{\partial^\alpha u(x_i, t_j)}{\partial_+ x^\alpha} &= \frac{1}{(\Delta x)^\alpha} \sum_{k=0}^{i+1} g_k u(x_i - (k-1)\Delta x, t_j) \\ &= \frac{1}{(\Delta x)^\alpha} \sum_{k=0}^{i+1} g_k u_{i-k+1, j} \end{aligned} \quad (3)$$

and

$$\frac{\partial^\alpha u(x_i, t_j)}{\partial_x^\alpha} = \frac{1}{(\Delta x)^\alpha} \sum_{k=0}^{n-i+1} g_k u(x_i + (k-1)\Delta x, t_j)$$

$$= \frac{1}{(\Delta x)^\alpha} \sum_{k=0}^{n-i+1} g_k u_{i+k-1, j}$$
(4)

Where  $g_0 = 1$  and  $g_k = (-1)^k \frac{\alpha(\alpha-1)\dots(\alpha-k+1)}{k!}$ ,  $k = 1, 2, \dots$

To improve the explicit finite difference method for solving the initial-boundary value problem (1) - (2), we substitute  $t = t_j$  in

equation (1) and replace the partial derivative  $\frac{\partial^2 u}{\partial t^2}$  with its central difference approximation to get:

$$\frac{u_{i,j+1} - 2u_{i,j} + u_{i,j-1}}{(\Delta t)^2} = c_{i,j} \frac{\partial^\alpha u_{i,j}}{\partial x^\alpha} + s_{i,j}$$
(5)

Where  $t_j = j\Delta t$ ,  $j = 0, 1, \dots, m$  and  $m$  is the number of subdivision of the interval  $[0, T]$ ,  $t \in R$ . Next, substitute equation (3) in equation (5) to obtain:

$$\frac{u_{i,j+1} - 2u_{i,j} + u_{i,j-1}}{(\Delta t)^2} = \frac{c_{i,j}}{(\Delta x)^\alpha} \sum_{k=0}^{i+1} g_k u_{i-k+1, j} + s_{i,j}$$

$i = 1, 2, \dots, n$ ;  $j = 0, 1, \dots, m-1$

(6)

Also, the initial and boundary conditions given by equation (2) become:

$$\left. \begin{aligned} u_{r,0} = u(x_i, 0) = f(x_i), \quad \frac{\partial u(x_i, 0)}{\partial t} = h(x_i) \quad \text{for } i = 0, 1, \dots, n \\ u_{0,j} = u(L, t_j) = 0, \quad u_{n,j} = u(R, t) = 0, \quad \text{for } j = 0, 1, \dots, m \end{aligned} \right\}$$

Now, using the central difference approximation to the initial derivative conditions, we can get:

$$\frac{1}{2\Delta t} (u_{i,1} - u_{i,-1}) = h(x_i), \quad i = 0, 1, \dots, n$$

Which can be written as:

$$u_{i,1} = u_{i,-1} + 2\Delta t h(x_i), \quad i = 0, 1, \dots, n$$

Furthermore, equation (6) becomes:

$$u_{i,j+1} = 2u_{i,j} - u_{i,j-1} + \frac{(\Delta t)^2 c_{i,j}}{(\Delta x)^\alpha} \sum_{k=0}^{i+1} g_k u_{i-k+1, j} + s_{i,j}(\Delta t)^2; \quad i = 1, 2, \dots, n; \quad j = 0, 1, \dots, m-1$$
(7)

Therefore:

$$u_{i,1} = 2u_{i,0} - u_{i,-1} + \frac{(\Delta t)^2 c_{i,0}}{(\Delta x)^\alpha} \sum_{k=0}^{i+1} g_k u_{i-k+1, 0} + s_{i,0}(\Delta t)^2$$
(8)

By substituting  $u_{i,-1} = u_{i,1} - 2\Delta t h(x_i)$  back into equation (8), one can show that  $u_{i,1}$  can be calculated from the following equation:

$$u_{i,1} = f_i + \frac{(\Delta t)^2 c_{i,0}}{2(\Delta x)^\alpha} \sum_{k=0}^{i+1} g_k f_{i-k+1, 0} + \frac{(\Delta t)^2}{2} s_{i,0} + \Delta t g_i; \quad i = 1, 2, \dots, n-1.$$

By evaluating the above equation for each  $i = 1, 2, \dots, n-1$ , one can get the values of  $u_{i,1}$ .

Then by evaluating equation (7) at  $i = 1, 2, \dots, n-1$  and  $j = 2, 3, \dots, m-1$ , one can get the numerical solution of equation (1). The resulting equation can be explicitly solved to give:

$$u_{i,j+1} - 2u_{i,j} + u_{i,j-1} = r \sum_{w=0}^{i+1} g_w u_{i-w+1, j}$$
(9)

Where  $r = \frac{k^2}{h^\alpha}$ .

It's well known that in the explicit finite difference treatment the PDE or the FPDE are replaced by an algebraic system of equations which can be written as the form

$$A\bar{u} = \bar{f},$$
(10)

Where,  $A$  is nonsingular coefficients matrix,  $\bar{u} = [u_{j,1}, u_{j,2}, \dots, u_{j,N-1}]^T$  and  $\bar{f} = [f_{j,1}, f_{j,2}, \dots, f_{j,N-1}]^T$ ,  $j = 1, 2, \dots, N-1$ .

Now, from the linear system of equations (10) which formed when fractional hyperbolic partial differential equation is solved by the explicit finite difference method, matrix  $A$  can be write as  $A = D - L - U$  where  $D$  is diagonal matrix  $A$ ,  $-L$  is strictly lower triangular parts of  $A$  and  $-U$  is strictly upper triangular parts of  $A$ . A preconditioner  $(I - kL)$ , where:  $1 \leq k < 2$  is used to modify the original system (10) to

$$(I - kL)A\bar{u} = (I - kL)\bar{f}$$
(11)

The resulted system of (11) called preconditioned explicit finite difference method.

### 3. Stability analysis

We can discuss the stability of the resulting equation (9) as the following:

$$\text{Let } g_0 = 1 \text{ and } g_w = (-1)^w \frac{\alpha(\alpha-1)k(\alpha-w-1)}{w}; \quad w = 1, 2, \dots; \quad 1 \leq \alpha \leq 2.$$

Hence,  $g \geq 0$ , for all values of  $i$ . Therefore:

$$\sum_{w=0}^{i+1} g_w \leq g_1 = -(-\alpha) = \alpha$$
(12)

The difference between the analytical and numerical solutions of the difference equation remains bounded as  $j$  increases.

Let the error  $E_{ij} = u(h_i, k_j) - u_{i,j}$  then the finite difference equation (9) is stable. Now, we have to find the stability condition under which the error  $E_{ij}$  is bounded. Smith [15] shows that the error  $E_{ij}$  can be written as:

$$E_{ij} = e^{\sqrt{-1}\beta i h} \zeta_j$$
(13)

By substituting equations (12), (13) into (9), we can get:

$$\varepsilon - 2 - \varepsilon^{-1} - r\alpha e^{\sqrt{-1}\beta h(1-w)} \leq 0$$
(14)

Assume that  $\lambda = \beta h(1-w)$  and substitute it into equation (14), it can be easily getting the following equation for  $R$  as:

$$\varepsilon^2 - (2 - r\alpha e^{\sqrt{-1}\lambda})\varepsilon + 1 = 0 \tag{15}$$

Let  $k = 2 + r\alpha e^{\sqrt{-1}\lambda}$  and  $|e^{\sqrt{-1}\lambda}| \leq 1$ . Hence the values of  $\varepsilon$  are:

$$\varepsilon_1 = \frac{k + \sqrt{k^2 - 4}}{2}, \quad \varepsilon_2 = \frac{k - \sqrt{k^2 - 4}}{2}$$

From equation (13), the error will not grow with time if

$$|e^{\sqrt{-1}\lambda}| \leq 1 \text{ for all real } \beta \tag{16}$$

Equation (16) is called the Von-Neumann's condition for stability. Therefore, we will use this equation to find the stability condition of the finite difference equation (9).

We can see that for  $r, \alpha$  and  $\beta$  are real, the stability will be given while  $\varepsilon_2$  gives instability. For  $-1 \leq k \leq 1$ , the only useful inequality is  $k \leq 1$  and then  $2 + r\alpha e^{\sqrt{-1}\lambda} \leq 1$ , where  $|e^{\sqrt{-1}\lambda}| \leq 1$ . There-

fore,  $r \leq \frac{-1}{\alpha}$  where  $1 \leq \alpha \leq 2$ .

Hence,  $|r| \leq \frac{1}{2}$  which is the stability condition. We can conclude

that the stability results in the finite partial differential equation case as generalization for the corresponding result in the classical hyperbolic partial differential equation.

By the same manner, we can observe that the preconditioned system (11) has the same stability condition because it has the same structure of (10).

### 4. Numerical results and discussion

Several numerical experiments have been conducted to show the superiority of the proposed method for solving fractional hyperbolic partial differential equation. To check the effectiveness of the proposed iterative method; we will use the following modal problem which is the fractional order partial differential equation [16]:

$$\frac{\partial^2 u}{\partial t^2} = \frac{1}{\Gamma(0.5)} \sqrt{x} \frac{\partial^{1.5} u}{\partial x^{1.5}} - 4x^2 + 2x^3 - (2.546)x^2 t^2 + (2.546)x t^2, 0 \leq x \leq 2, 0 \leq t \leq 1,$$

Together with initial and zero Dirichlet boundary conditions:

$$\left. \begin{aligned} u(x, 0) = 0, \quad u_t(x, 0) = 0, \quad 0 \leq x \leq 2 \\ u(0, t) = 0, \quad u(1, t) = 0, \quad 0 \leq t \leq 1 \end{aligned} \right\}$$

This problem has the exact solution  $u(x, t) = x^2(x - 2)t^2$ . Firstly, the numerical solutions based on preconditioned explicit finite difference and the original explicit finite difference were compared with the exact solution to show the accuracy of these iterative methods. Figure 1 shows that the accuracy of the mentioned methods is acceptable. Furthermore, we can observe that the preconditioned explicit finite difference have better accuracy than the original one.

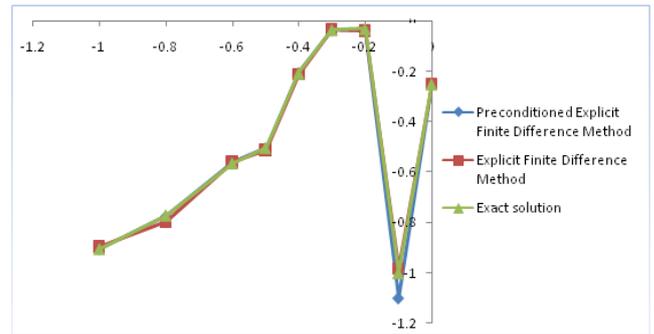


Fig. 1: Comparison of Exact Solution to the Numerical Solution from the Preconditioned and the Original Explicit Finite Difference

In addition to that, Comparisons for Number of iterations and Elapsed time of proposed method and the original one are made for the particular mesh size. Table 1. shows that the proposed method has the lowest number of iterations and elapsed time.

Table 1: Comparison of Number of Iterations and Elapsed Time of Proposed and Original Explicit Finite Difference Methods for Solving the Model Problem

N	Explicit Finite Difference method			Preconditioned Explicit Finite Difference method		
	Number of iterations	Elapse d time (sec.)	Average absolute Error	Number of iterations	Elapse d time (sec.)	Average absolute Error
74	1572	0.0974	0.0244	1438	0.0183	0.0192
96	1946	0.1733	0.0194	1873	0.1532	0.0134
162	2511	0.2521	0.0131	2360	0.2403	0.0114
186	2815	0.2844	0.0094	2670	0.2501	0.0063
222	3178	0.3423	0.0064	2981	0.2904	0.0043
246	3397	0.2874	0.0024	3165	0.1863	0.0015

### 5. Conclusion

In this paper, we have formulated new preconditioned explicit finite difference iterative method for solving fractional hyperbolic partial differential equations. From observation of all experimental results, it can be concluded that the proposed scheme may be a good alternative to solve this type of fractional differential equation and many other numerical problems.

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