

Multiplication theorems for multi-variable and multi-index Bessel functions

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Abstract

In this article, we derive multiplication theorems for the 3-variable 2-parameter Bessel functions $J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2)$ and 2-index 5-variable 3-parameter Bessel functions $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3)$ using the generating function method. Further, we derive multiplication theorems for functions related to $J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2)$ and $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3)$. Furthermore, we establish a multiplication theorem for N-index Bessel functions $J_{m_1, m_2, \dots, m_N}(\lambda x)$.

Keywords: Multiplication theorem, Generalized Bessel functions, Double-index Bessel functions.

1 Introduction

Theory of special functions plays an important role in the formalism of mathematical physics. Bessel functions (BF) [13], are among the most important special functions, with very diverse applications to physics, engineering and mathematical physics ranging from abstract number theory and theoretical astronomy to concrete problems of physics and engineering.

Dattoli and his co-workers introduced and discussed various generalizations of BF within purely mathematical and applicative contexts (see for example [1-5,7,9-11]), as the success of any generalized special functions depend on its usefulness in applications and on the existence of effective numerical procedures for its computation. Generalized Bessel functions (GBF) have become a powerful tool to investigate the dynamical aspects of physical problems such as electron scattering by an intense linearly polarized laser wave, multi-photon processes and undulator radiation. The analytical and numerical study of GBF has revealed their interesting properties, which in some sense can be regarded as an extension of the properties of BF to a 2-dimensional domain. In this connection, the relevance of GBF and their multi-variable extension in mathematical physics has been emphasized, since they provide analytical solutions to partial differential equations such as the multi-dimensional diffusion equation, the Schrödinger and Klein-Gordon equations. A useful complement to the theory of GBF is offered by the introduction of 3-variable 2-parameter Bessel functions (3V2PBF) $J_n(x, y, z; \tau_1, \tau_2)$ defined as ([8], p.222)

$$J_n(x, y, z; \tau_1, \tau_2) = \sum_{l=-\infty}^{\infty} \tau_2^l J_{n-3l}(x, y; \tau_1) J_l(z), \quad 0 < |\tau_1|, |\tau_2| < \infty, \quad (1)$$

where τ_1, τ_2 are arbitrary complex parameters. The generating function for 3V2PBF is given as

$$\sum_{n=-\infty}^{\infty} J_n(x, y, z; \tau_1, \tau_2) t^n = \exp \left(\frac{x}{2} \left(t - \frac{1}{t} \right) + \frac{y}{2} \left(t^2 \tau_1 - \frac{1}{t^2 \tau_1} \right) + \frac{z}{2} \left(t^3 \tau_2 - \frac{1}{t^3 \tau_2} \right) \right). \quad (2)$$

In particular, we note that

$$J_n(x, y, z; 1, 1) = J_n(x, y, z), \quad (3)$$

where $J_n(x, y, z)$ denotes the 3-variable generalized Bessel functions (3VGBF) defined by the generating function [2]

$$\sum_{n=-\infty}^{\infty} J_n(x, y, z) t^n = \exp \left(\frac{x}{2} \left(t - \frac{1}{t} \right) + \frac{y}{2} \left(t^2 - \frac{1}{t^2} \right) + \frac{z}{2} \left(t^3 - \frac{1}{t^3} \right) \right). \quad (4)$$

Further, we note the following link

$$J_n(x, y, 0; \tau_1 \rightarrow \tau, 0) = J_n(x, y; \tau), \quad (5)$$

where $J_n(x, y; \tau)$ denotes the 2-variable 1-parameter generalized Bessel functions (2V1PGBF) defined by the generating function [1]

$$\sum_{n=-\infty}^{\infty} J_n(x, y; \tau) t^n = \exp \left(\frac{x}{2} \left(t - \frac{1}{t} \right) + \frac{y}{2} \left(t^2 \tau - \frac{1}{t^2 \tau} \right) \right) \quad (6)$$

and specified by the series

$$J_n(x, y; \tau) = \sum_{l=-\infty}^{\infty} \tau_2^l J_{n-2l}(x) J_l(y). \quad (7)$$

Also, we note that

$$J_n(x, y, 0; 1, 1) = J_n(x, y), \quad (8)$$

where $J_n(x, y)$ denotes the 2-variable generalized Bessel functions (2VGBF) defined by the generating function [6]

$$\sum_{n=-\infty}^{\infty} J_n(x, y) t^n = \exp \left(\frac{x}{2} \left(t - \frac{1}{t} \right) + \frac{y}{2} \left(t^2 - \frac{1}{t^2} \right) \right). \quad (9)$$

Moreover, it is evident that

$$J_n(x, 0, 0; 1, 1) = J_n(x), \quad (10)$$

where $J_n(x)$ denotes the ordinary Bessel functions (BF) defined by the generating function [13]

$$\sum_{n=-\infty}^{\infty} J_n(x) t^n = \exp \left(\frac{x}{2} \left(t - \frac{1}{t} \right) \right). \quad (11)$$

We considered the 2-index 5-variable 3-parameter Bessel functions (2I5V3PBF) $J_{m,n}(x, y, z, w, h; \tau_1, \tau_2, \tau_3)$ defined as

$$J_{m,n}(x, y, z, w, h; \tau_1, \tau_2, \tau_3) = \sum_{l=-\infty}^{\infty} \tau_3^l J_{m-l}(x, z; \tau_1) J_{n-l}(y, w; \tau_2) J_l(h), \quad (12)$$

with the following generating function

$$\begin{aligned} \sum_{m,n=-\infty}^{\infty} J_{m,n}(x, y, z, w, h; \tau_1, \tau_2, \tau_3) u^m v^n &= \exp \left(\frac{x}{2} \left(u - \frac{1}{u} \right) + \frac{y}{2} \left(v - \frac{1}{v} \right) + \frac{z}{2} \left(u^2 \tau_1 - \frac{1}{u^2 \tau_1} \right) \right. \\ &\quad \left. + \frac{w}{2} \left(v^2 \tau_2 - \frac{1}{v^2 \tau_2} \right) + \frac{h}{2} \left(u v \tau_3 - \frac{1}{u v \tau_3} \right) \right). \end{aligned} \quad (13)$$

In particular, we note that

$$J_{m,n}(x, y, z, w, h; 1, 1, 1) = J_{m,n}(x, y, z, w, h), \quad (14)$$

where $J_{m,n}(x, y, z, w, h)$ denotes the 2-index 5-variable Bessel functions (2I5VBF) defined by the generating function ([6], p.188(3.23))

$$\sum_{m,n=-\infty}^{\infty} J_{m,n}(x, y, z, w, h) u^m v^n = \exp \left(\frac{x}{2} \left(u - \frac{1}{u} \right) + \frac{y}{2} \left(v - \frac{1}{v} \right) + \frac{z}{2} \left(u^2 - \frac{1}{u^2} \right) + \frac{w}{2} \left(v^2 - \frac{1}{v^2} \right) + \frac{h}{2} \left(uv - \frac{1}{uv} \right) \right). \quad (15)$$

Further, we note the following link

$$J_{m,n}(x, y, 0, 0, h \rightarrow z; \tau_1, \tau_2, \tau_3 \rightarrow \tau) = J_{m,n}(x, y, z; \tau), \quad (16)$$

where $J_{m,n}(x, y, z; \tau)$ denotes the 2-index 3-variable 1-parameter Bessel functions (2I3V1PBF) defined by the generating function ([10], p.344(2))

$$\sum_{m,n=-\infty}^{\infty} J_{m,n}(x, y, z; \tau) u^m v^n = \exp \left(\frac{x}{2} \left(u - \frac{1}{u} \right) + \frac{y}{2} \left(v - \frac{1}{v} \right) + \frac{z}{2} \left(uv\tau - \frac{1}{uv\tau} \right) \right). \quad (17)$$

Furthermore, it is evident that

$$J_{m,n}(x, y, 0, 0, h \rightarrow z; \tau_1, \tau_2, 1) = J_{m,n}(x, y, z), \quad (18)$$

where $J_{m,n}(x, y, z)$ denotes the 2-index 3-variable Bessel functions (2I3VBF) defined by the generating function ([7], p.3639(13))

$$\sum_{m,n=-\infty}^{\infty} J_{m,n}(x, y, z) u^m v^n = \exp \left(\frac{x}{2} \left(u - \frac{1}{u} \right) + \frac{y}{2} \left(v - \frac{1}{v} \right) + \frac{z}{2} \left(uv - \frac{1}{uv} \right) \right). \quad (19)$$

Also, we note that

$$J_{m,n}(x, y \rightarrow x, 0, 0, h \rightarrow x; \tau_1, \tau_2, 1) = J_{m,n}(x), \quad (20)$$

where $J_{m,n}(x)$ denotes the 2-index Bessel functions (2IBF) defined by the generating function ([7], p.3637(1))

$$\sum_{m,n=-\infty}^{\infty} J_{m,n}(x) u^m v^n = \exp \left(\frac{x}{2} \left(\left(u - \frac{1}{u} \right) + \left(v - \frac{1}{v} \right) + \left(uv - \frac{1}{uv} \right) \right) \right). \quad (21)$$

Also, the 2-index 1-parameter Bessel functions (2I1PBF) $J_{m,n}(x; \tau)$ defined by means of the series [7]

$$J_{m,n}(x; \tau) = \sum_{l=-\infty}^{\infty} \tau^l J_{m-l}(x) J_{n-l}(x) J_l(x). \quad (22)$$

Moreover, we note that

$$J_{m,n}(x, 0, z \rightarrow y, 0, 0; 1, \tau_2, \tau_3) = J_m(x, y), \quad (23)$$

or

$$J_{m,n}(0, y \rightarrow x, 0, w \rightarrow y, 0; \tau_1, 1, \tau_3) = J_n(x, y) \quad (24)$$

and

$$J_{m,n}(x, 0, 0, 0, 0; \tau_1, \tau_2, \tau_3) = J_m(x), \quad (25)$$

or

$$J_{m,n}(0, y \rightarrow x, 0, 0, 0; \tau_1, \tau_2, \tau_3) = J_n(x) \quad (26)$$

Next, we recall the 3-variable generalized Hermite polynomials (3VGHP) $H_n(x, y, z)$ defined by the generating function ([3], p.511(17))

$$\sum_{n=0}^{\infty} H_n(x, y, z) \frac{t^n}{n!} = \exp(xt + yt^2 + zt^3), \quad (27)$$

which for $z = 0$ reduce to the 2VHKdFP $H_n(x, y)$ defined by the generating function [5]

$$\sum_{n=0}^{\infty} H_n(x, y) \frac{t^n}{n!} = \exp(xt + yt^2) \quad (28)$$

and specified by the series

$$H_n(x, y) = n! \sum_{m=0}^{\lfloor \frac{n}{2} \rfloor} \frac{(-y)^m (2x)^{n-2m}}{m!(n-2m)!}. \quad (29)$$

The 2-index 5-variable generalized Hermite polynomials (2I5VGHP) $H_{m,n}(x, y, z, w, h)$ defined by the generating function [6]

$$\sum_{m,n=0}^{\infty} H_{m,n}(x, y, z, w, h) \frac{u^m v^n}{m! n!} = \exp(xu + yv + zu^2 + wv^2 + hvw), \quad (30)$$

which for $w = h = 0$ reduce to the 2-index 3-variable Hermite polynomials (2I3VHP) $H_{m,n}(x, y, z)$ defined by [14]

$$H_{m,n}(x, y, z) = \sum_{p=0}^{\min(m,n)} \frac{(\frac{x}{2})^{m-p} (\frac{y}{2})^{n-p} (\frac{z}{2})^p m! n!}{p! (m-p)! (n-p)!}. \quad (31)$$

Further, the 2-index Hermite polynomials (2IHP) $H_{m,n}(x)$ are specified by the series [5]

$$H_{m,n}(x) = \sum_{p=0}^{\min(m,n)} \frac{(\frac{x}{2})^{m+n-p} m! n!}{p! (m-p)! (n-p)!}. \quad (32)$$

Hermite polynomials are useful complements to the theory of generalized Bessel functions. The link between these two families can be easily expressed as follows:

$$J_n(x, y, z; 1, 1) = \sum_{r=0}^{\infty} \frac{1}{r!(n+r)!} H_{n+r}\left(\frac{x}{2}, \frac{y}{2}, \frac{z}{2}\right) H_r\left(\frac{-x}{2}, \frac{-y}{2}, \frac{-z}{2}\right) \quad (33)$$

and

$$J_{m,n}(x, y, z, w, h; 1, 1, 1) = \sum_{r,s=0}^{\infty} \frac{1}{r! s! (m+r)! (n+s)!} H_{m+r, n+s}\left(\frac{x}{2}, \frac{y}{2}, \frac{z}{2}, \frac{w}{2}, \frac{h}{2}\right) H_{r,s}\left(\frac{-x}{2}, \frac{-y}{2}, \frac{-z}{2}, \frac{-w}{2}, \frac{-h}{2}\right). \quad (34)$$

Theory of generalized Bessel functions has been developed mainly for their practical importance in many physical and engineering problems, ranging from non-dipolar scattering to surface diffraction theory. Motivated by the importance of multivariable forms of Bessel functions and their link with other forms of generalized Bessel functions, in this paper we establish multiplication theorems for the 3V2PBF $J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2)$ and 2I5V3PBF $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3)$ using the generating function method. Also, we derive multiplication theorems for functions related to $J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2)$ and $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3)$ as applications. Finally, we prove a multiplication theorem for N-index Bessel functions (N-IBF) $J_{m_1, m_2, \dots, m_N}(\lambda x)$ as a possible extension of our method.

In section 2, we derive multiplication theorems for the 3V2PBF $J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2)$ 2I5V3PBF $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3)$. In section 3, we consider applications of our main results. In section 4, concluding remarks are given.

2 Multiplication theorems for 3V2PBF $J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2)$ and 2I5V3PBF $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3)$

In this section, we prove the following multiplication theorems involving 3V2PBF $J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2)$ and 2I5V3PBF $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3)$:

Theorem 2.1 *The following multiplication theorem for the 3V2PBF $J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2)$ holds true:*

$$J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2) = \lambda^{-n} \sum_{m=0}^{\infty} \frac{1}{m!} J_{n-m}\left(x, y, z; \frac{\lambda^2 \tau_1}{\mu}, \frac{\lambda^3 \tau_2}{\nu}\right) H_m\left(\frac{x}{2}(\lambda^2 - 1), \frac{\lambda^2 \tau_1 y}{2\mu}(\mu^2 - 1), \frac{\lambda^3 \tau_2 z}{2\nu}(\nu^2 - 1)\right). \quad (35)$$

Proof We start with the identity

$$\begin{aligned}
& \exp \left(\frac{\lambda x}{2} \left(t - \frac{1}{t} \right) + \frac{\mu y}{2} \left(t^2 \tau_1 - \frac{1}{t^2 \tau_1} \right) + \frac{\nu z}{2} \left(t^3 \tau_2 - \frac{1}{t^3 \tau_2} \right) \right) \\
&= \exp \left(\frac{x}{2} \left(\frac{t}{\lambda} - \frac{\lambda}{t} \right) \right) \exp \left(\frac{y}{2} \left(\frac{t^2 \tau_1}{\mu} - \frac{\mu}{t^2 \tau_1} \right) \right) \exp \left(\frac{z}{2} \left(\frac{t^3 \tau_2}{\nu} - \frac{\nu}{t^3 \tau_2} \right) \right) \\
& \exp \left(\frac{x}{2} (\lambda^2 - 1) \left(\frac{t}{\lambda} \right) + \frac{\lambda^2 \tau_1 y}{2\mu} (\mu^2 - 1) \left(\frac{t}{\lambda} \right)^2 + \frac{\lambda^3 \tau_2 z}{2\nu} (\nu^2 - 1) \left(\frac{t}{\lambda} \right)^3 \right), \tag{36}
\end{aligned}$$

which on using the definitions (11) and (27), on the r.h.s. becomes

$$\begin{aligned}
& \exp \left(\frac{\lambda x}{2} \left(t - \frac{1}{t} \right) + \frac{\mu y}{2} \left(t^2 \tau_1 - \frac{1}{t^2 \tau_1} \right) + \frac{\nu z}{2} \left(t^3 \tau_2 - \frac{1}{t^3 \tau_2} \right) \right) \\
&= \sum_{h,k,p=-\infty}^{\infty} \frac{1}{\lambda^h} J_h(x) \left(\frac{\tau_1}{\mu} \right)^k \left(\frac{\tau_2}{\nu} \right)^p t^{h+2k+3p} J_k(y) J_p(z) \sum_{m=0}^{\infty} \left(\frac{t}{\lambda} \right)^m \frac{1}{m!} \\
& H_m \left(\frac{x}{2} (\lambda^2 - 1), \frac{\lambda^2 \tau_1 y}{2\mu} (\mu^2 - 1), \frac{\lambda^3 \tau_2 z}{2\nu} (\nu^2 - 1) \right). \tag{37}
\end{aligned}$$

Now, replacing h by $s - 2k$ and using the relation (7) in the r.h.s. of Eq.(37), we find

$$\begin{aligned}
& \exp \left(\frac{\lambda x}{2} \left(t - \frac{1}{t} \right) + \frac{\mu y}{2} \left(t^2 \tau_1 - \frac{1}{t^2 \tau_1} \right) + \frac{\nu z}{2} \left(t^3 \tau_2 - \frac{1}{t^3 \tau_2} \right) \right) \\
&= \sum_{s,p=-\infty}^{\infty} \frac{1}{\lambda^s} J_s \left(x, y; \frac{\lambda^2 \tau_1}{\mu} \right) \left(\frac{\tau_2}{\nu} \right)^p t^{s+3p} J_p(z) \sum_{m=0}^{\infty} \left(\frac{t}{\lambda} \right)^m \frac{1}{m!} \\
& H_m \left(\frac{x}{2} (\lambda^2 - 1), \frac{\lambda^2 \tau_1 y}{2\mu} (\mu^2 - 1), \frac{\lambda^3 \tau_2 z}{2\nu} (\nu^2 - 1) \right). \tag{38}
\end{aligned}$$

Again, replacing s by $q - 3p$ and using the relation (1) in the r.h.s. of Eq.(38), we get

$$\begin{aligned}
& \exp \left(\frac{\lambda x}{2} \left(t - \frac{1}{t} \right) + \frac{\mu y}{2} \left(t^2 \tau_1 - \frac{1}{t^2 \tau_1} \right) + \frac{\nu z}{2} \left(t^3 \tau_2 - \frac{1}{t^3 \tau_2} \right) \right) \\
&= \sum_{q=-\infty}^{\infty} \sum_{m=0}^{\infty} \left(\frac{t}{\lambda} \right)^{q+m} \frac{1}{m!} J_q \left(x, y, z; \frac{\lambda^2 \tau_1}{\mu}, \frac{\lambda^3 \tau_2}{\nu} \right) \\
&\quad H_m \left(\frac{x}{2} (\lambda^2 - 1), \frac{\lambda^2 \tau_1 y}{2\mu} (\mu^2 - 1), \frac{\lambda^3 \tau_2 z}{2\nu} (\nu^2 - 1) \right). \tag{39}
\end{aligned}$$

Further, replacing q by $n - m$ in the r.h.s. and using definition (2) on the l.h.s. of above equation, we find

$$\begin{aligned}
\sum_{n=-\infty}^{\infty} J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2) t^n &= \sum_{n=-\infty}^{\infty} t^n \lambda^{-n} \sum_{m=0}^{\infty} \frac{1}{m!} J_{n-m} \left(x, y, z; \frac{\lambda^2 \tau_1}{\mu}, \frac{\lambda^3 \tau_2}{\nu} \right) \\
&\quad H_m \left(\frac{x}{2} (\lambda^2 - 1), \frac{\lambda^2 \tau_1 y}{2\mu} (\mu^2 - 1), \frac{\lambda^3 \tau_2 z}{2\nu} (\nu^2 - 1) \right). \tag{40}
\end{aligned}$$

Finally, on equating the coefficients of like powers of t in Eq. (40), we get the assertion (35) of Theorem 2.1.

Theorem 2.2 *The following multiplication theorem for the 2I5V3PBF $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3)$ holds true:*

$$\begin{aligned}
J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3) &= \lambda^{-m} \mu^{-n} \sum_{i,j=0}^{\infty} \frac{\lambda^i \mu^j}{i! j!} J_{m-i, n-j} \left(x, y, z, w, h; \frac{\lambda^2 \tau_1}{\nu}, \frac{\mu^2 \tau_2}{\eta}, \frac{\lambda \mu \tau_3}{\beta} \right) \\
&\quad H_{i,j} \left(\frac{x}{2\lambda} (\lambda^2 - 1), \frac{y}{2\mu} (\mu^2 - 1), \frac{z \tau_1}{2\nu} (\nu^2 - 1), \frac{w \tau_2}{2\eta} (\eta^2 - 1), \frac{h \tau_3}{2\beta} (\beta^2 - 1) \right). \tag{41}
\end{aligned}$$

Proof We start with the identity

$$\begin{aligned}
& \exp \left(\frac{\lambda x}{2} \left(u - \frac{1}{u} \right) + \frac{\mu y}{2} \left(v - \frac{1}{v} \right) + \frac{\nu z}{2} \left(u^2 \tau_1 - \frac{1}{u^2 \tau_1} \right) + \frac{\eta w}{2} \left(v^2 \tau_2 - \frac{1}{v^2 \tau_2} \right) + \frac{\beta h}{2} \left(uv \tau_3 - \frac{1}{uv \tau_3} \right) \right) \\
&= \exp \left(\frac{x}{2} \left(\frac{u}{\lambda} - \frac{\lambda}{u} \right) \right) \exp \left(\frac{y}{2} \left(\frac{v}{\mu} - \frac{\mu}{v} \right) \right) \exp \left(\frac{z}{2} \left(\frac{u^2 \tau_1}{\nu} - \frac{\nu}{u^2 \tau_1} \right) \right) \exp \left(\frac{w}{2} \left(\frac{v^2 \tau_2}{\eta} - \frac{\eta}{v^2 \tau_2} \right) \right) \\
&\quad \exp \left(\frac{h}{2} \left(\frac{uv \tau_3}{\beta} - \frac{\beta}{uv \tau_3} \right) \right) \exp \left(\frac{x}{2} \left(\frac{\lambda^2 - 1}{\lambda} \right) u + \right. \\
&\quad \left. \frac{y}{2} \left(\frac{\mu^2 - 1}{\mu} \right) v + \frac{z \tau_1}{2} \left(\frac{\nu^2 - 1}{\nu} \right) u^2 + \frac{w \tau_2}{2} \left(\frac{\eta^2 - 1}{\eta} \right) v^2 + \frac{h \tau_3}{2} \left(\frac{\beta^2 - 1}{\beta} \right) uv \right), \tag{42}
\end{aligned}$$

which on using the definitions (11) and (30), on the r.h.s. becomes

$$\begin{aligned}
& \exp \left(\frac{\lambda x}{2} \left(u - \frac{1}{u} \right) + \frac{\mu y}{2} \left(v - \frac{1}{v} \right) + \frac{\nu z}{2} \left(u^2 \tau_1 - \frac{1}{u^2 \tau_1} \right) + \frac{\eta w}{2} \left(v^2 \tau_2 - \frac{1}{v^2 \tau_2} \right) + \frac{\beta h}{2} \right. \\
&\quad \left. \left(uv \tau_3 - \frac{1}{uv \tau_3} \right) \right) = \sum_{p,q,r,s,t=-\infty}^{\infty} \frac{u^{p+2r+t} v^{q+2s+t}}{\lambda^h \mu^q \nu^r \eta^s \beta^t} \tau_1^r \tau_2^s \tau_3^t J_p(x) J_q(y) J_r(z) J_s(w) J_t(h) \\
&\quad \sum_{i,j=0}^{\infty} \frac{u^i v^j}{i! j!} H_{i,j} \left(\frac{x}{2\lambda} (\lambda^2 - 1), \frac{y}{2\mu} (\mu^2 - 1), \frac{z \tau_1}{2\nu} (\nu^2 - 1), \frac{w \tau_2}{2\eta} (\eta^2 - 1), \frac{h \tau_3}{2\beta} (\beta^2 - 1) \right). \tag{43}
\end{aligned}$$

Now, replacing p by $l - 2r$ and q by $k - 2s$ and using the relation (7) in the r.h.s. of Eq.(43), we find

$$\begin{aligned}
& \exp \left(\frac{\lambda x}{2} \left(u - \frac{1}{u} \right) + \frac{\mu y}{2} \left(v - \frac{1}{v} \right) + \frac{\nu z}{2} \left(u^2 \tau_1 - \frac{1}{u^2 \tau_1} \right) + \frac{\eta w}{2} \left(v^2 \tau_2 - \frac{1}{v^2 \tau_2} \right) + \frac{\beta h}{2} \right. \\
&\quad \left. \left(uv \tau_3 - \frac{1}{uv \tau_3} \right) \right) = \sum_{l,k,t=-\infty}^{\infty} \frac{u^{l+t} v^{k+t}}{\lambda^l \mu^k \beta^t} \left(\frac{\lambda \mu \tau_3}{\beta} \right)^t J_l \left(x, z; \frac{\lambda^2 \tau_1}{\nu} \right) J_k \left(y, w; \frac{\mu^2 \tau_2}{\eta} \right) J_t(h)
\end{aligned}$$

$$\sum_{i,j=0}^{\infty} \frac{u^i v^j}{i! j!} H_{i,j} \left(\frac{x}{2\lambda}(\lambda^2 - 1), \frac{y}{2\mu}(\mu^2 - 1), \frac{z\tau_1}{2\nu}(\nu^2 - 1), \frac{w\tau_2}{2\eta}(\eta^2 - 1), \frac{h\tau_3}{2\beta}(\beta^2 - 1) \right). \quad (44)$$

Again, replacing l by $\theta - t$ and k by $\phi - t$ and using the relation (12) in the r.h.s. of Eq.(44), we get

$$\begin{aligned} & \exp \left(\frac{\lambda x}{2} \left(u - \frac{1}{u} \right) + \frac{\mu y}{2} \left(v - \frac{1}{v} \right) + \frac{\nu z}{2} \left(u^2 \tau_1 - \frac{1}{u^2 \tau_1} \right) + \frac{\eta w}{2} \left(v^2 \tau_2 - \frac{1}{v^2 \tau_2} \right) + \frac{\beta h}{2} \right. \\ & \left. \left(uv \tau_3 - \frac{1}{uv \tau_3} \right) \right) = \sum_{\theta, \phi=-\infty}^{\infty} \sum_{i,j=0}^{\infty} \frac{u^{\theta+i} v^{\phi+j}}{\lambda^{\theta} \mu^{\phi} i! j!} J_{\theta, \phi} \left(x, y, z, w, h; \frac{\lambda^2 \tau_1}{\nu}, \frac{\mu^2 \tau_2}{\eta}, \frac{\lambda \mu \tau_3}{\beta} \right) \\ & H_{i,j} \left(\frac{x}{2\lambda}(\lambda^2 - 1), \frac{y}{2\mu}(\mu^2 - 1), \frac{z\tau_1}{2\nu}(\nu^2 - 1), \frac{w\tau_2}{2\eta}(\eta^2 - 1), \frac{h\tau_3}{2\beta}(\beta^2 - 1) \right). \end{aligned} \quad (45)$$

Further, replacing θ by $m - i$ and ϕ by $n - j$ in the r.h.s. and using the definition (13) on the l.h.s. of above equation, we find

$$\begin{aligned} & \sum_{m,n=-\infty}^{\infty} J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3) u^m v^n = \sum_{m,n=-\infty}^{\infty} u^m v^n \lambda^{-m} \mu^{-n} \\ & \sum_{i,j=0}^{\infty} \frac{\lambda^i \mu^j}{i! j!} J_{m-i, n-j} \left(x, y, z, w, h; \frac{\lambda^2 \tau_1}{\nu}, \frac{\mu^2 \tau_2}{\eta}, \frac{\lambda \mu \tau_3}{\beta} \right) \\ & H_{i,j} \left(\frac{x}{2\lambda}(\lambda^2 - 1), \frac{y}{2\mu}(\mu^2 - 1), \frac{z\tau_1}{2\nu}(\nu^2 - 1), \frac{w\tau_2}{2\eta}(\eta^2 - 1), \frac{h\tau_3}{2\beta}(\beta^2 - 1) \right). \end{aligned} \quad (46)$$

Finally, on equating the coefficients of like powers of u and v in Eq. (46), we get the assertion (41) of Theorem 2.2.

3 Applications

We discuss some applications of the multiplication theorems obtained in the preceding section.

I. Taking $z = \tau_2 = 0$ and replacing τ_1 by τ in Eq. (35) and using Eq. (5), we obtain the multiplication theorem for 2V1PGBF $J_n(\lambda x, \mu y; \tau)$ ([12],p.145(10))

$$J_n(\lambda x, \mu y; \tau) = \lambda^{-n} \sum_{m=0}^{\infty} \frac{1}{m!} J_{n-m} \left(x, y; \frac{\lambda^2 \tau}{\mu} \right) H_m \left(\frac{x}{2}(\lambda^2 - 1), \frac{\lambda^2 \tau y}{2\mu}(\mu^2 - 1) \right), \quad (47)$$

where $H_n(x, y)$ denotes the 2VHKdFP given by Eq. (28).

II. Replacing τ_1 by $\frac{\mu}{\lambda^2}$ and τ_2 by $\frac{\nu}{\lambda^3}$ in Eq. (35) and using Eq. (33), we get

$$\begin{aligned} J_n \left(\lambda x, \mu y, \nu z; \frac{\mu}{\lambda^2}, \frac{\nu}{\lambda^3} \right) &= \lambda^{-n} \sum_{m,r=0}^{\infty} \frac{1}{m! r! (n+r)!} H_{n-m+r} \left(\frac{x}{2}, \frac{y}{2}, \frac{z}{2} \right) H_r \left(\frac{-x}{2}, \frac{-y}{2}, \frac{-z}{2} \right) \\ & H_m \left(\frac{x}{2}(\lambda^2 - 1), \frac{y}{2}(\mu^2 - 1), \frac{z}{2}(\nu^2 - 1) \right), \end{aligned} \quad (48)$$

where $J_n(x, y, z; \tau_1, \tau_2)$ denotes the 3V2PBF given by Eq. (2) and $H_n(x, y, z)$ denotes the 3VGHP given by Eq. (27).

III. Taking $\tau_1 = \tau_2 = 1$ in Eq. (35) and using Eq. (3), we obtain the multiplication theorem for 3VGBF $J_n(\lambda x, \mu y, \nu z)$

$$J_n(\lambda x, \mu y, \nu z) = \lambda^{-n} \sum_{m=0}^{\infty} \frac{1}{m!} J_{n-m} \left(x, y, z; \frac{\lambda^2}{\mu}, \frac{\lambda^3}{\nu} \right) H_m \left(\frac{x}{2}(\lambda^2 - 1), \frac{\lambda^2 y}{2\mu}(\mu^2 - 1), \frac{\lambda^3 z}{2\nu}(\nu^2 - 1) \right). \quad (49)$$

where $J_n(x, y, z)$ denotes the 3VGBF given by Eq. (4).

IV. Taking $\tau_1 = \tau_2 = \tau_3 = 1$ in Eq. (41) and using Eq. (14), we get the multiplication theorem for 2I5VBF $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h)$

$$\begin{aligned} J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h) &= \lambda^{-m} \mu^{-n} \sum_{i,j=0}^{\infty} \frac{\lambda^i \mu^j}{i! j!} J_{m-i, n-j} \left(x, y, z, w, h; \frac{\lambda^2}{\nu}, \frac{\mu^2}{\eta}, \frac{\lambda \mu}{\beta} \right) \\ &\quad H_{i,j} \left(\frac{x}{2\lambda} (\lambda^2 - 1), \frac{y}{2\mu} (\mu^2 - 1), \frac{z}{2\nu} (\nu^2 - 1), \frac{w}{2\eta} (\eta^2 - 1), \frac{h}{2\beta} (\beta^2 - 1) \right). \end{aligned} \quad (50)$$

where $J_{m,n}(x, y, z, w, h; \tau_1, \tau_2, \tau_3)$ denotes the 2I5V3PBF given by Eq. (13) and $H_{i,j}(x, y, z, w, h)$ denotes the 2I5VGHP given by Eq. (30).

V. Taking $z = w = 0$, $\tau_3 = 1$ and replacing h by z in Eq. (41) and using Eq. (18), we find the multiplication theorem for 2I3VBF $J_{m,n}(\lambda x, \mu y, \beta z)$

$$J_{m,n}(\lambda x, \mu y, \beta z) = \lambda^{-m} \mu^{-n} \sum_{i,j=0}^{\infty} \frac{\lambda^i \mu^j}{i! j!} J_{m-i, n-j} \left(x, y, z; \frac{\lambda \mu}{\beta} \right) H_{i,j} \left(\frac{x}{2\lambda} (\lambda^2 - 1), \frac{y}{2\mu} (\mu^2 - 1), \frac{z}{2\beta} (\beta^2 - 1) \right), \quad (51)$$

where $J_{m,n}(x, y, z; \tau)$ denotes the 2I3V1PBF given by Eq. (17) and $H_{i,j}(x, y, z)$ denotes the 2I3VGHP given by Eq. (31). Further, replacing y and z by x and μ and β by λ in Eq. (51) and using Eq. (20), we obtain a result ([12],p.146(11)).

VI. Replacing τ_1 by $\frac{\nu}{\lambda^2}$, τ_2 by $\frac{\eta}{\mu^2}$ and τ_3 by $\frac{\beta}{\lambda \mu}$ in Eq. (41) and using Eq. (34), we get

$$\begin{aligned} J_{m,n} \left(\lambda x, \mu y, \nu z, \eta w, \beta h; \frac{\nu}{\lambda^2}, \frac{\eta}{\mu^2}, \frac{\beta}{\lambda \mu} \right) &= \lambda^{-m} \mu^{-n} \sum_{i,j,r,s=0}^{\infty} \frac{\lambda^i \mu^j}{i! j! r! s! (m+r)! (n+s)!} \\ &\quad H_{m-i+r, n-j+s} \left(\frac{x}{2}, \frac{y}{2}, \frac{z}{2}, \frac{w}{2}, \frac{h}{2} \right) H_{r,s} \left(\frac{-x}{2}, \frac{-y}{2}, \frac{-z}{2}, \frac{-w}{2}, \frac{-h}{2} \right) \\ &\quad H_{i,j} \left(\frac{x}{2\lambda} (\lambda^2 - 1), \frac{y}{2\mu} (\mu^2 - 1), \frac{z\tau_1}{2\nu} (\nu^2 - 1), \frac{w\tau_2}{2\eta} (\eta^2 - 1), \frac{h\tau_3}{2\beta} (\beta^2 - 1) \right). \end{aligned} \quad (52)$$

VII. Taking $y = w = h = 0$, $\tau_1 = 1$ and replacing z by y in Eq. (41) and using Eq. (23), we get the multiplication theorem for 2VBF $J_m(\lambda x, \nu y)$

$$J_m(\lambda x, \nu y) = \lambda^{-m} \sum_{i=0}^{\infty} \frac{\lambda^i}{i!} J_{m-i} \left(x, y; \frac{\lambda^2}{\nu} \right) H_i \left(\frac{x}{2\lambda} (\lambda^2 - 1), \frac{y}{2\nu} (\nu^2 - 1) \right), \quad (53)$$

or, taking $x = z = h = 0$, $\tau_2 = 1$ and replacing y by x and w by y in Eq. (41) and using Eq. (24), we get

$$J_n(\mu x, \eta y) = \mu^{-n} \sum_{j=0}^{\infty} \frac{\mu^j}{j!} J_{n-j} \left(x, y; \frac{\mu^2}{\eta} \right) H_j \left(\frac{x}{2\mu} (\mu^2 - 1), \frac{y}{2\eta} (\eta^2 - 1) \right), \quad (54)$$

where $J_n(x, y; \tau)$ denotes the 2V1PBF given by Eq. (6) and $J_n(x, y)$ denotes the 2VGBF given by Eq. (9).

VIII. Taking $y = z = w = h = 0$ in Eq. (41) and using Eqs. (25) and (29), we find the multiplication theorem for BF $J_m(\lambda x)$ ([12],p.150(25))

$$J_m(\lambda x) = \lambda^{-m} \sum_{i=0}^{\infty} \frac{1}{i!} (x(\lambda^2 - 1))^i J_{m-i}(x), \quad (55)$$

or, taking $x = z = w = h = 0$ and replacing y by x in Eq. (41) and using Eqs. (26) and (29), we get

$$J_n(\mu x) = \mu^{-n} \sum_{j=0}^{\infty} \frac{1}{j!} (x(\mu^2 - 1))^j J_{n-j}(x), \quad (56)$$

where $J_n(x)$ denotes the ordinary Bessel functions.

4 Conclusion

In section 2, we have derive multiplication theorems for the 3V2PBF $J_n(\lambda x, \mu y, \nu z; \tau_1, \tau_2)$ and 2I5V3PBF $J_{m,n}(\lambda x, \mu y, \nu z, \eta w, \beta h; \tau_1, \tau_2, \tau_3)$ using the generating function method. We remark that this process can be extended to establish multiplication theorem for the N-IBF $J_{m_1, m_2, \dots, m_N}(\lambda x)$. First, we recall the N-IBF $J_{m_1, m_2, \dots, m_N}(x)$ defined by the generating function [6]

$$\sum_{m_j=-\infty}^{\infty} J_{m_1, \dots, m_N}(x) \prod_{j=1}^N u_j^{m_j} = \exp \left(\frac{x}{2} \left(\sum_{j=1}^N \left(u_j - \frac{1}{u_j} \right) + \left(\prod_{j=1}^N u_j - \frac{1}{\prod_{j=1}^N u_j} \right) \right) \right), \quad (57)$$

with the following series

$$J_{m_1, \dots, m_N}(x) = \sum_{l=-\infty}^{\infty} \left(\prod_{j=1}^N J_{m_j-l}(x) \right) J_l(x). \quad (58)$$

Also, we can define the N-index 1-variable 1-parameter Bessel functions (N-I1V1PBF) $J_{m_1, \dots, m_N}(x; \tau)$ by means of the series

$$J_{m_1, \dots, m_N}(x; \tau) = \sum_{l=-\infty}^{\infty} \tau^{(N-1)l} \left(\prod_{j=1}^N J_{m_j-l}(x) \right) J_l(x). \quad (59)$$

Now, we start with the identity

$$\begin{aligned} \exp \left(\frac{\lambda x}{2} \left(\sum_{j=1}^N \left(u_j - \frac{1}{u_j} \right) + \left(\prod_{j=1}^N u_j - \frac{1}{\prod_{j=1}^N u_j} \right) \right) \right) &= \exp \left(\frac{x}{2} \left(\sum_{j=1}^N \left(\frac{u_j}{\lambda} - \frac{\lambda}{u_j} \right) \right) \right) \\ &\quad \exp \left(\frac{x}{2} \left(\prod_{j=1}^N \frac{u_j}{\lambda} - \prod_{j=1}^N \frac{\lambda}{u_j} \right) \right) \exp \left(\frac{x}{2} \left(\frac{\lambda^2 - 1}{\lambda} \right) \left(\sum_{j=1}^N u_j + \prod_{j=1}^N u_j \right) \right), \end{aligned} \quad (60)$$

which on using the definition (11), on the r.h.s. becomes

$$\begin{aligned} &\exp \left(\frac{\lambda x}{2} \left(\sum_{j=1}^N \left(u_j - \frac{1}{u_j} \right) + \left(\prod_{j=1}^N u_j - \frac{1}{\prod_{j=1}^N u_j} \right) \right) \right) \\ &= \sum_{p_j, r_j, q=-\infty}^{\infty} \frac{\prod_{j=1}^N u_j^{p_j+r_j+q}}{\lambda^{p_j+q} r_1!, r_2!, \dots, r_N!} J_{p_1}(x) J_{p_2}(x) \dots J_{p_N}(x) J_q(x) H_{r_1, r_2, \dots, r_N} \left(x \left(\frac{\lambda^2 - 1}{\lambda} \right) \right), \end{aligned} \quad (61)$$

where

$$H_{r_1, r_2, \dots, r_N}(x) = \sum_{k=0}^{\min(r_1, r_2, \dots, r_N)} \left(\frac{x}{2} \right)^{\sum_{j=1}^N r_j - (N-1)k} \frac{r_1!, r_2!, r_3!, \dots, r_N!}{k!(r_1-k)!(r_2-k)!\dots(r_N-k)!}. \quad (62)$$

Now, replacing p_j by $s_j - q$ and using the relation (59) in the r.h.s. of Eq.(61), we find

$$\begin{aligned} \exp \left(\frac{\lambda x}{2} \left(\sum_{j=1}^N \left(u_j - \frac{1}{u_j} \right) + \left(\prod_{j=1}^N u_j - \frac{1}{\prod_{j=1}^N u_j} \right) \right) \right) &= \sum_{s_j, r_j=-\infty}^{\infty} \frac{\prod_{j=1}^N u_j^{s_j+r_j}}{\lambda^{s_i} r_1!, r_2!, \dots, r_N!} J_{s_1, s_2, \dots, s_N}(x; \lambda) \\ &\quad H_{r_1, r_2, \dots, r_N} \left(x \left(\frac{\lambda^2 - 1}{\lambda} \right) \right). \end{aligned} \quad (63)$$

Further, replacing s_j by $m_j - r_j$ in the r.h.s. and using definition (57) on the l.h.s. of above equation, we find

$$\begin{aligned} \sum_{m_j=-\infty}^{\infty} J_{m_1, \dots, m_N}(\lambda x) \prod_{j=1}^N u_j^{m_j} &= \sum_{m_j, r_j=-\infty}^{\infty} \frac{\prod_{j=1}^N u_j^{m_j}}{\lambda^{m_j-r_j} r_1!, r_2!, \dots, r_N!} J_{m_1-r_1, m_2-r_2, \dots, m_N-r_N}(x; \lambda) \\ &\quad H_{r_1, r_2, \dots, r_N} \left(x \left(\frac{\lambda^2 - 1}{\lambda} \right) \right). \end{aligned} \quad (64)$$

Finally, on equating the coefficients of like powers of $\prod_{j=1}^N u_j^{m_j}$ in Eq. (63), we get multiplication theorem for the N-IBF $J_{m_1, m_2, \dots, m_N}(\lambda x)$ as follows

$$J_{m_1, \dots, m_N}(\lambda x) = \lambda^{-m_1-m_2-\dots-m_N} \sum_{r_j=-\infty}^{\infty} \frac{\lambda^{r_1+r_2+\dots+r_N}}{r_1!, r_2!, \dots, r_N!} J_{m_1-r_1, m_2-r_2, \dots, m_N-r_N}(x; \lambda) H_{r_1, r_2, \dots, r_N} \left(x \left(\frac{\lambda^2 - 1}{\lambda} \right) \right). \quad (65)$$

The theory of BF is rich and wide and certainly provides an inexhaustible field of research. A large number of functions are recognized as belonging to the BF family. The above multiplication theorem can be viewed as a multi-index extension of multiplication theorem (3.9) for the BF $J_m(x)$. Exploring the possibility of using the method outlined here to derive the multiplication theorems for other special functions is a further research problems.

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